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Please cite as:

Antonio Pasculli, Monia Calista, Nicola Sciarra (2018), "Variability of local stress states resulting from the application of Monte Carlo and finite difference methods to the stability study of a selected slope"

Engineering Geology, Vol. 245, pp. 370-389, doi: <https://doi.org/10.1016/j.enggeo.2018.09.009>

1 **Variability of local stress states resulting from the application of Monte Carlo and finite**
2 **difference methods to the stability study of a selected slope**

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10

11 **Abstract**

12

13 In slope stability analysis it is decisive to represent, in the most realistic way, the heterogeneity of the
14 soil characteristics. The Monte Carlo technique is one of the suitable mathematical tools that could
15 be useful for this purpose. The result of this kind of approach is a set of multiple possible realizations
16 of the same system, characterized by different spatial distribution of the numerical values of the
17 geotechnical parameters. The numerical stress states outcomes, obtained through a stability analysis,
18 would depend on the models employed to define the spatial parameters distribution and, further, on
19 the utilized geomechanical models. Accordingly, first of all, we selected a suitable mathematical
20 model, based on the Monte Carlo technique. Then, we integrated it into a commonly used FDM
21 (Finite Difference Method) commercial code (FLAC2D). Finally, we studied the stability of an actual
22 slope, considered as a test case (Lettomanoppello, Chieti, Italy), whose geological and geomechanical
23 characteristics make it suitable for the application of the selected code. In this step, in order to simplify
24 the analyses, we considered only dry conditions. Therefore, we discussed how the stochastic
25 distributions of the local stress states, resulting from hundreds of repeated runs, were correlated to
26 the mechanical parameters distributions obtained by the Monte Carlo application and assigned to each
27 points as inputs. The results of the test application suggest that our approach can be used to identify
28 sectors of the slopes more sensitive to the variability of the input values of the main geotechnical
29 parameters, which would require a more accurate modelling and monitoring.

30

31 Keywords: numerical modeling; geological modeling; random parameters distribution; Monte Carlo
32 method; statistics.

33

34 **1. Introduction**

35

36 A suitable and effective mathematical tool, capable to include the complexity of soil structures,
37 caused also by weathering (Esposito et al. 2013; Pasculli et al., 2017) may be provided by statistical
38 approaches. Some authors, including Vanmarcke (1977), Fenton and Vanmarcke (1990), Rabczuk et
39 al. (2007), Andrade et al. (2008), Bossi et al. (2016), Vu-Bac et al. (2016), highlighted the importance
40 of considering the possible numerical values of the geotechnical parameters as random variables,
41 belonging to an assumed Probability Distribution Function (PDF), for example a Gaussian Function,
42 particularly in the case of the study of the stability of granular media. Of course, the obtained results
43 were based on algorithms that can calculate random numerical values. Current computing techniques
44 are not able to generate truly random numbers. Many strategies were proposed, based, for instance,
45 on the internal clock of the computer (Banks et al., 2005). However, all the generated sets of
46 numerical values have a cyclicity. Accordingly, after a certain number of generations, even if very
47 long (10^{36} as well), the obtained sequences repeat themselves. For this reason the values are called
48 ‘pseudo-random numbers’. The PDF can be acquired as the outcome of multiple runs, all with the
49 same inputs, but with different parameter values, computed automatically by a selected mathematical
50 algorithm based on pseudo-random generation, integrated as intrinsic routine into most codes. This
51 kind of approach is called ‘Monte Carlo’, already applied in the field of slope stability analyses
52 (among many others: Sundell et al., 2016). In the context of the study of slope stability, many
53 geomechanical and mathematical-numerical approaches have been proposed in the international
54 literature. Among these, Rabczuk and Areias (2006) and Rabczuk et al. (2007) proposed interesting
55 approaches for modeling shear bands with cohesive geological materials, through the important mesh-
56 free numerical technique, which avoids the cumbersome construction of meshes, like *Smoothed*
57 *Particle Hydrodynamics* (SPH). This kind of approach was long used in many fields including those
58 where it is necessary to simulate large deformations or transition from solid-like to fluid like of the
59 materials or in the study of debris flow (among many others, Minatti and Pasculli 2010, 2011, Pasculli
60 et al. 2013, 2014). However, despite these techniques are very valid, the most common commercial
61 computer codes such as those belonging to the FLAC family, widely and commonly used in
62 professional practice, are based on different approaches. Consequently, for the purpose of this paper,
63 we have selected this type of computer code, in its 2D version. FLAC2D is based on Finite Difference
64 Method (FDM) applied to continuum mechanics (among many others: Jiang et al. 2001; Sitharam et

65 al. 2007; Calista et al. 2015a, 2015b, 2016;). The continuum is discretized by a grid as suitable as
66 possible for the system under study, faults and sliding surfaces included. A significant feature of the
67 FLAC2D code is the option for assigning the numerical values of the geotechnical parameters or
68 suitable constitutive equations to each grid cell, according to specific rules selected by the user,
69 through the interactive language called FISH (based on the C compiler). At this stage of the research,
70 for the sake of simplicity, the connectivity between the value of mechanical parameters assigned to a
71 point and the values of near others points was not implemented, observing that the common
72 experimental connectivity scale (less than 1 meter, Vanmarke, 1977; Ching and Wang, 2016) was
73 lower than the selected dimension of the average distance among points (ranging from 10 to 20
74 meters) resulting from the selected discretization of the continuum. Some applications to actual slope
75 systems, using both FLAC2D and FLAC3D, can already be found in Calista et al. (2015a and b),
76 Pasculli and Sciarra (2006a), Pasculli et al., (2006b), Sciarra et al., (2006), where the first attempts to
77 exploit Monte Carlo techniques not only for 2D, but also for 3D simulations, were pursued. However,
78 the works available in the framework of the international literature, among which Vanmarke (1977),
79 Fenton and Vanmarke (1990) and more recently Bossi et al. (2016), Ching and Wang, (2016), Sundell
80 et al. (2016), have not sufficiently deepened yet the important issue of how common numerical
81 approaches, exploited by the most used commercial computer, may influence and modify the spatial
82 propagation of the uncertainties. Accordingly, for the sake of simplicity, we proposed an approach
83 considering a surrogate deterministic dependence of all geotechnical variables on just one parameter:
84 the Unit volume weight. The comparisons between the variability of the inputs and outputs data,
85 resulting from the application of FLAC, were carried out, according to the scope of the paper, by the
86 commonly utilized commercial code STATGRAPHICS Centurion XVI, 2010. Available options
87 were reported in Table A.1, 'Exploited statistical tools', and in Table A.2, 'Fitting curve' (Cox and
88 Hinkley, 1974, Casella and Berger, 2001). Then, after this type of analysis, for particular identified
89 sectors of the system under study, an enhanced modelling, like multivariate approach, may be
90 required as a successive step. Recently, Vu-Bac et al. (2016) and Hamdia et al. (2017) have proposed
91 an interesting approach aimed at realizing a global sensitivity analysis of the variables that affect the
92 selected system in order to quantify the significance of the input parameters and their joint effects on
93 the output by varying all the parameters at the same time. In these two papers the polynomial chaos
94 expansions (Ghanem and Spanos,1991) were applied. However, Ghanem and Spanos (1991) also
95 exploited the Monte Carlo method in order to test the analytical chaos expansions. Thus, in our paper,
96 we found reasonable to select a model based on the Monte Carlo method (Pasculli and Sciarra (2002,
97 2003) as an effective alternative to chaos function expansion. Moreover a Yielding Index (F_{YI}),
98 similar to a local safety factor, on which the global factor may depend, calculated in each mesh of the

99 numerically discretized continuum, was introduced as an indicative parameter of the local stress state
100 and of the possible local risk of rupture. Usually, the uncertainties related to the determination of
101 important physical and geotechnical parameters can be introduced by assigning a PDF of these
102 parameters to each node of the discretized spatial domain. Accordingly, the calculated values of
103 parameters that depend on these quantities, such as for example the local stability factors, will be
104 affected by some variability and uncertainties, described by the resulting PDFs, not necessarily
105 similar to those provided for the input parameters. Therefore, the comparison between the input PDFs
106 and the resulting calculated PDFs, can provide, as discussed below, the indications, at least for the
107 system under study, on how the uncertainties introduced in the inputs propagated through the entire
108 spatial domain. After a brief discussion of the mathematical model, details of the algorithm
109 implementation are given. The interactive FISH code was used to make the FLAC2D more suitable
110 for Monte Carlo multiple realizations of possible real structural configurations. For the purpose of
111 this study, rainfall infiltration and filtration were not included. Accordingly, the selected slope was
112 considered in dry conditions. Then, the results related to 650 Monte Carlo realizations of the spatial
113 assignment of mechanical parameters values, belonging to a prior selected Gaussian numerical
114 ensemble, based on real data and representative of an actual slope test case, are discussed. The choice
115 of the number of realizations was conditioned by only the request to limit the calculation time. Finally,
116 the statistical distributions of the cohesion coefficient related to the representative area of the slope,
117 obtained through the Monte Carlo approach, were compared to the resulting distributions of the
118 calculated F_{YI} , applying a commonly used commercial code (STATGRAPHICS).

119 120 **2. Mathematical algorithms and code implementation description**

121
122 For the sake of clarity, a brief sketch of the implemented mathematical model, already partially
123 discussed in our previous works (Pasculli. and Sciarra 2003, 2006, for instance), is presented and
124 more insights are given. Each mechanical parameter was assumed to be a random variable Ψ_{rand} ,
125 and accordingly, affected by some uncertainties. However, a sufficient number of samples, necessary
126 to build a suitable statistics, is not always available, as in the case of the test selected for this work.
127 On the other hand, an extensive bibliography is also dedicated to the description of the range of
128 variability of the most important parameters characterizing the materials that constitute soils.
129 Therefore, even for the purposes of the paper, we considered synthetic data based, however, on
130 experimental values with the related variability drawn from the bibliography and reported in Table 1,
131 regarding material resulting from the geophysical surveys carried out for the test case. In particular,
132 the ‘Unit volume weight’ of the Clay material (γ_{rand}) was assumed to be a random variable Ψ_{rand}

133 affected, from Table 1, by the following uncertainties: $19.5 \text{ (kN/m}^3) \leq \gamma_{rand} \leq 21.0 \text{ (kN/m}^3)$.
 134 Another important statistical parameter was the Standard deviation σ_Ψ that, following the adopted
 135 approach, had to be deduced from the data shown in the same Table 1. We estimated this parameter
 136 on the base of the well *Chebyshev inequality*, according to which about 93% of the numerical value
 137 of Ψ_{rand} lies in the range: $\Psi_{mean} - 4\sigma_\Psi \leq \Psi_{rand} \leq \Psi_{mean} + 4\sigma_\Psi$, where $\Psi_{mean} = (\Psi_{min} +$
 138 $\Psi_{max})/2$ was the estimated mean, (Clay: $\Psi_{mean} = (19.5 + 21.0)/2 = 20.25 \text{ kN/m}^3$),
 139 independently from the probability distribution function to which the random variable belongs. By
 140 consequence, we assumed that in general: $\Psi_{min} = \Psi_{mean} - 4\sigma_\Psi$, $\Psi_{max} = \Psi_{mean} + 4\sigma_\Psi$.
 141 Moreover the resulting standard deviation was estimated in the following way:

$$142 \quad \sigma_\Psi = \frac{(\Psi_{max} - \Psi_{min})}{8} = 0.125 \cdot (\Psi_{max} - \Psi_{min}) \quad (1)$$

143
 144 In particular, for the Clay: $\sigma_\Psi = \frac{(\Psi_{max} - \Psi_{min})}{8} = 0.125 \cdot (\Psi_{max} - \Psi_{min}) = 0.125 \cdot (21.0 - 19.5) =$
 145 0.0625 kN/m^3 .

146 Following the above procedure, we were able to define the generic random variable, assumed suitable
 147 for simulating the numerical values of the mechanical parameters characterizing the soil under study:
 148
 149

$$150 \quad \Psi_{rand} = \Psi_{mean} + (\sigma_\Psi \cdot g_{norm}) \quad (2)$$

151 where:

$$152 \quad g_{norm} = \sqrt{[-2 \cdot \ln(y_{ran1})]} \cdot \cos(2\pi \cdot y_{ran2}) \quad (3)$$

153
 154 was the Standard Gaussian Distribution (SGD) provided by the algorithm proposed by Box and
 155 Muller (1958), while y_{ran1} and y_{ran2} were two non-correlated (pseudo) random variables
 156 uniformly distributed in the interval (0, 1) (provided directly by the FLAC's intrinsic function
 157 URAND). For the Clay: $\Psi_{rand}(\text{Unit volume weight}) \equiv \gamma_{rand} = [20.25 + (0.0625 \cdot$
 158 $g_{norm})] \text{ kN/m}^3$. Truncated probability distribution (Hesse et al., 2014) was not considered here,
 159 thus, by consequence, the SGD, and accordingly Eq. (3), may provide random values outside of the
 160 assumed range. Hence, in this case, the value was set equal to the mean. Furthermore, again for the
 161 sake of simplicity, we assumed that random spatial variations of each mechanical parameters were in
 162 direct accordance with the random spatial variation of the 'Unit volume weight' of the selected
 163 material (Deposits, Clay, Clay marl and Chalk alteration, Limestone) through a simple linear law:
 164
 165

166

$$167 \quad \Psi_{rand}(x, y) = (\Psi_{max} - \Psi_{min}) \frac{[\gamma_{rand}(x, y) - \gamma_{min}]}{(\gamma_{max} - \gamma_{min})} + \Psi_{min} \quad (4)$$

168

169 where $\Psi_{rand}(x, y)$ and $\gamma_{rand}(x, y)$ were the stochastic numerical value of the selected parameter
170 (friction, cohesion, elastic modulus) and the ‘Unit volume weight’ assigned to the point P(x,y),
171 through the application of the proposed algorithm. Furthermore, an important underlying assumption
172 was that the used statistical distribution was not spatially stationary. Now, through the above
173 approach, the values of the parameter were estimated on only the basis of a pure application of
174 statistics, without considering the alteration due to the specific conditions and interactions to which
175 each sample would be affected in situ. Accordingly, a combination of a deterministic and random
176 terms due, respectively, to the lithostatic pressure or other loads at macroscopic scale and to a local
177 grains contact distribution at small scale had to be included.

178 Then the deviation of the in situ $\gamma_{rand_insitu}(x, y)$ and the estimated value $\gamma_{rand_lab}(x, y) = \gamma_{mean} +$
179 $(\sigma_{\gamma_lab} \cdot g_norm1)$, to be measured in a ‘virtual laboratory’, resulting from the previous
180 consideration in accordance with Eqs. (2) and (3), was assumed to be as much higher as the difference
181 between the value that the selected parameter would have had if there had been only the deterministic
182 term due to lithological deterministic load $\mu_{\gamma}(x, y)$ and the intrinsic random value of that parameter,
183 $\gamma_{rand_lab}(x, y)$, multiplied by a random factor ε :

184

$$185 \quad \gamma_{rand_insitu}(x, y) - \gamma_{rand_lab}(x, y) = [\mu_{\gamma}(x, y) - \gamma_{rand_lab}(x, y)] \cdot \varepsilon \quad (5)$$

186

187 The ε parameter was a random factor that ‘weighed’ the deviation between the deterministic trend
188 and the random value of the parameter associated to the selected point P(x,y). The random factor was
189 assumed to be a normal Gaussian term (with unitary mean and standard deviation) similar to Eq. (1)
190 in which $(\Psi_{max} - \Psi_{min})$ was equal to one: $\varepsilon = 1. + 0.125 \cdot g_norm2$. It should be noted that
191 g_norm1 and g_norm2 were two different realizations of the SGD (Eq. 3), thus four non correlated
192 (pseudo) random variables were requested. The deterministic term $\mu_{\gamma}(x, y)$ was assumed to be equal
193 to the following relation:

194

$$195 \quad \mu_{\gamma}(x, y) = (\gamma_{max} - \gamma_{min}) \frac{[h(x, y) - h_{min}(y)]}{[h_{max}(y) - \gamma_{min}(y)]} + \gamma_{min} \quad (6)$$

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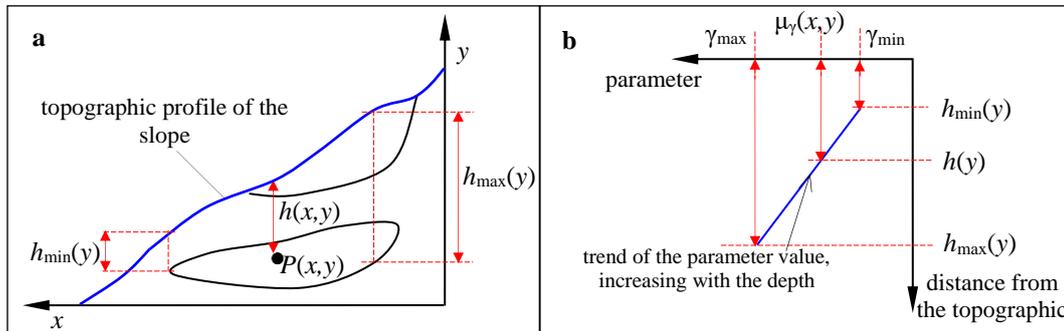
197 where: $h(x, y)$ was the depth of a generic point from the topographic level; $h_{min}(y)$ and $h_{max}(y)$
198 were, respectively, the minimum and the maximum depths at which a generic point embedded in a

199 specific sector, could be located (see lens in Fig. 1a). Fig. 1b shows the selected variation law of the
 200 parameters along depth. Finally, after simple manipulation, the following was the relation on which
 201 the implemented algorithms were based:

$$203 \gamma_{\text{insitu}}(x, y) = \mu_{\gamma}(x, y) + [\mu_{\gamma}(x, y) - \gamma_{\text{mean}}] \cdot 0.125 \cdot g_{\text{norm}2} - \sigma_{\gamma} \cdot 0.125 \cdot g_{\text{norm}1} \cdot g_{\text{norm}2} \quad (7)$$

204
 205 In Appendix we reported a flow chart (Fig. A.1.) of the adopted algorithm and a figure (Fig. A.2.),
 206 freely adapted from Pasculli and Sciarra (2003).

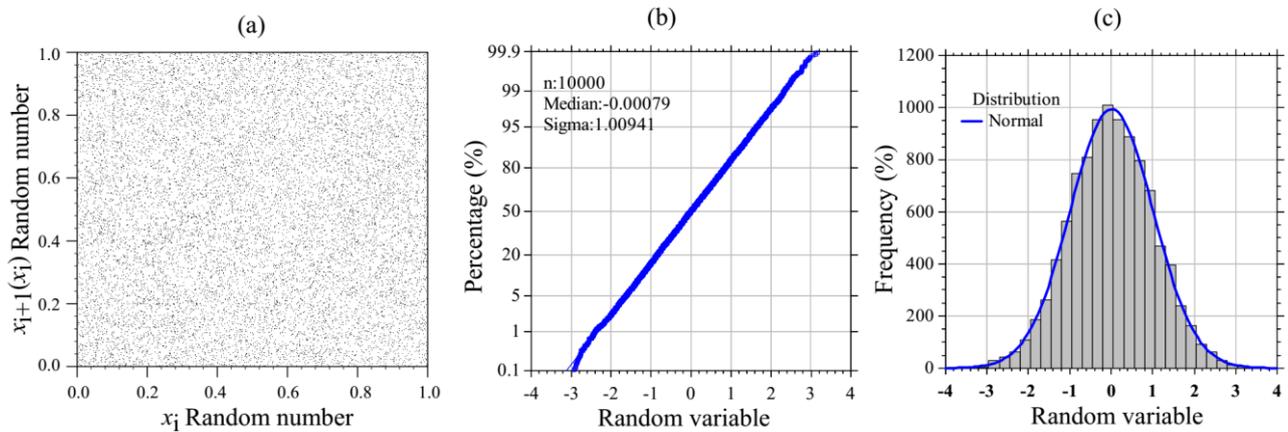
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 208 **Fig. 1.** Sketch of a generic, synthetic slope: a) geometrical parameters exploited in the mathematical
 209 model; b) selected trend of the parameters along depth.



217
 218 **2.1. Monte Carlo implementation test**

219
 220 The proposed Monte Carlo approach was based on the elaboration of (pseudo) random numbers,
 221 uniformly distributed within the range [0, 1], obtained essentially through both intrinsic algorithms
 222 included by default into the operative system and the used numerical tools (commercial such as FLAC
 223 or implemented by compilers such as Fortran, or C++). In this paper, in order to provide a different
 224 random number distribution, such as Gaussian distribution, the algorithm proposed by Box and
 225 Muller (1958), previously detailed, was selected. However, for its correct utilization, the (pseudo)
 226 random variables $y_{\text{ran}1}$ and $y_{\text{ran}2}$ (Eq. 3), uniformly distributed in the interval [0, 1] and
 227 calculated in sequence, should not be correlated. Hence, in Fig. 2a, the generated couples of 10,000
 228 pseudo random numbers, obtained by FLAC's intrinsic function URAND, are shown. In particular,
 229 along the abscissa and the ordinate axes, the i -th x random number corresponding to $y_{\text{ran}1}$ and the
 230 $(i+1)$ -th x random number corresponding to $y_{\text{ran}2}$, are respectively reported as black dot. The
 231 resulting dots distribution is clearly completely random (no trends at all). Thus, the absence of any
 232 correlation between x_i and successively generated x_{i+1} random number satisfies the requirement for
 233 the proper application of the Box Muller algorithm.

235 **Fig. 2.** a) Correlation test between x_i and successively generated x_{i+1} random numbers; b) Normality
 236 test; c) FLAC pseudo random numbers realization compared to Gaussian curve.
 237



238
 239
 240 Despite a correct application of both the Box-Muller's algorithm and the FLAC's intrinsic routine
 241 URAND, the realization of the distribution of random numbers, however, may suffer from an
 242 asymmetry around the mean. As shown in Fig. 2b and c, the random number distribution function,
 243 obtained according to the above detailed approaches, on which the Monte Carlo strategy adopted in
 244 this paper is based, is affected by a slight negative skewness (James, 2006).
 245 In short, the first step was to create a normal distribution of 10,000 random numbers (Fig. 2c).
 246 Accordingly, for each exploited mechanical parameter, an ensemble containing 10,000 possible
 247 random numerical values was generated through the expression (2), once for all the simulations. Then,
 248 at the beginning of each of the 650 realizations, the required numerical values input given by the
 249 expression (7), in particular the cohesion coefficients (function of the Unit volume weight, eq. 4) that
 250 should be assigned to each meshes, were randomly extracted from the associated ensemble by an
 251 automatic procedure. Thus, for each node, a different resulting statistical distribution of 650 F_{YI}
 252 numerical values was acquired. The cohesion coefficients were selected for further analyses and
 253 comparison.

254
 255 *2.2. FISH code implementation*

256
 257 In order to implement the algorithms briefly discussed above, we developed a FISH code, aimed at
 258 equipping the FLAC2D with the capability to apply Monte Carlo multiple realizations of possible
 259 actual structural configurations. In particular, Figs. A.3a and A.3b (in Appendix), show the FISH
 260 instructions, with some comments there reported, which were implemented in order to define,
 261 respectively, the distances from the topographic level of a selected point, as detailed in Fig. 1, and

262 the random value of the density assigned to it, following the algorithms characterized by Eqs. (2), (3),
 263 (4) and (5).

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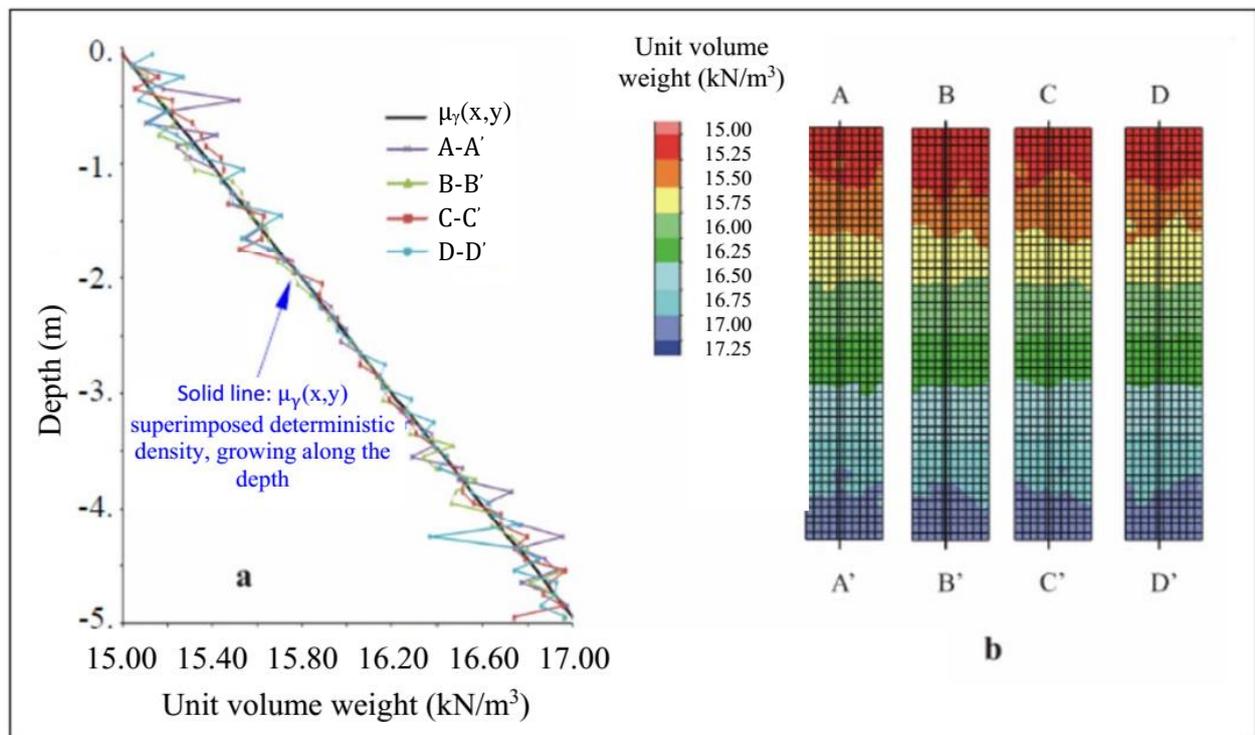
265 *2.2.1. FISH implementation test*

266

267 The algorithm, implemented into the FLAC2D, was first tested on a simple rectangular brick, using
 268 a 1 m wide and 5 m high grid, discretized with 0.1 m side square mesh. The aim of the proposed
 269 model is to attribute to each point geo-mechanical characteristics, affected by variability and, as a
 270 working hypothesis, by a linear trend with depth. Furthermore, each realization must be randomly
 271 different from the others. Accordingly, the Fig. 3b shows four different realizations with different
 272 random numerical values of the selected parameter (in this case the Unit volume weight), with a
 273 superimposed linear trend along the depth (Fig. 3a). The coloured contour plots of the four random
 274 realizations of the density and their graphs with the depth along a given vertical are reported. To this
 275 respect we can say that the results are ‘reasonable’ because they were the results we pursued to obtain.

276 **Fig. 3.** a) unit volume weight graphs in depth along the verticals A-A', B-B' and C-C'; b) contour
 277 plots of the unit volume weight for 4 realization.

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283 2.3. Yielding Index

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285 The selected code is based on the Finite Difference Method (FDM). Accordingly, the continuum is
286 divided in many meshes. In each mesh, the Mohr-Coulomb rupture criterion for all type of materials
287 was selected. The 2D stress state within each mesh can be expressed in terms of principal stress σ_1
288 and σ_3 . In the Mohr-Coulomb diagram, the associated stress state is identified by a circle (circle 'a'
289 in Fig. 4 freely adapted from Itasca Manual), characterized by a radius r_a . The element will reach the
290 rupture conditions when the circle will be tangent to the envelope of rupture (circle 'b' in Fig. 4)
291 defined, according to the Mohr-Coulomb criterion, through the well-known relation: $\tau = c' +$
292 $\sigma' \tan\varphi'$, where τ , c' , σ' and φ' are, respectively, the failure shear stress, the applied effective normal
293 stress, the effective angle of internal friction. In this paper we assumed that the values of the
294 mechanical parameters are referred to dry conditions. The ratio between the radii of the two circles
295 may be defined as 'Yielding Index', $F_{YI} = r_a/r_b$, somehow representing a local factor of safety. The
296 Yielding Index F_{YI} was selected as the key parameter for analyzing and collecting the results coming
297 from 650 different realizations.

298

299 **Fig. 4.** Mohr-Coulomb diagram (freely adapted from FLAC manual).

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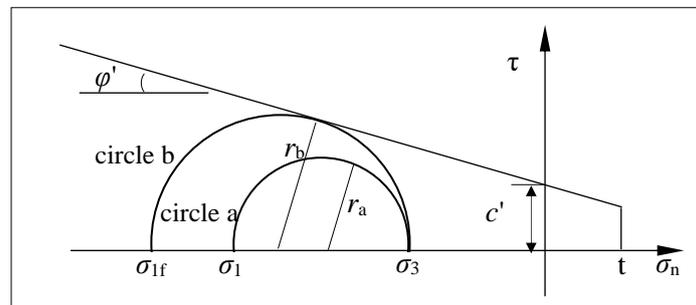
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313 **3. Test case**

314 The above discussed methodology was applied to an actual slope section. The study area, selected as
315 test case, is situated close to the town of Lettomanoppello (Pescara, Italy). It is located in the central-
316 eastern margin of the Apennines, along the north-western foothills of the relief of the Maiella (Fig.
317 5).

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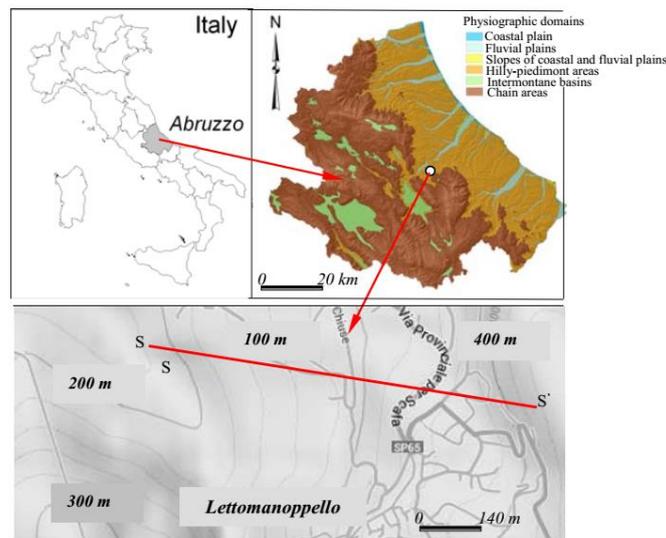
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325 **Fig. 5.** Geographic map of the investigated area



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328 The slope, on which the village of Lettomanoppello is located, is historically known for its
329 susceptibility to landslides. In fact, the literature sources indicate continuing movements since 1908.
330 The area includes Quaternary continental deposits (detritus deposits, limestone debris hetero-metric,
331 landslide body) and Mio-Pleistocene marine deposits (clay; clay, marl and chalk alternation;
332 limestone). The proposed approach was applied to a particularly significant section (section S-S' in
333 Fig. 5). In Fig. 6, the related outline of the section and the numerical FLAC grid, are reported. As
334 shown in Fig. 6a, the geological section is particularly complex because it features numerous sub
335 vertical faults and, in particular, a surface of discontinuity between Quaternary continental deposits
336 and Mio-Pleistocene marine deposits.

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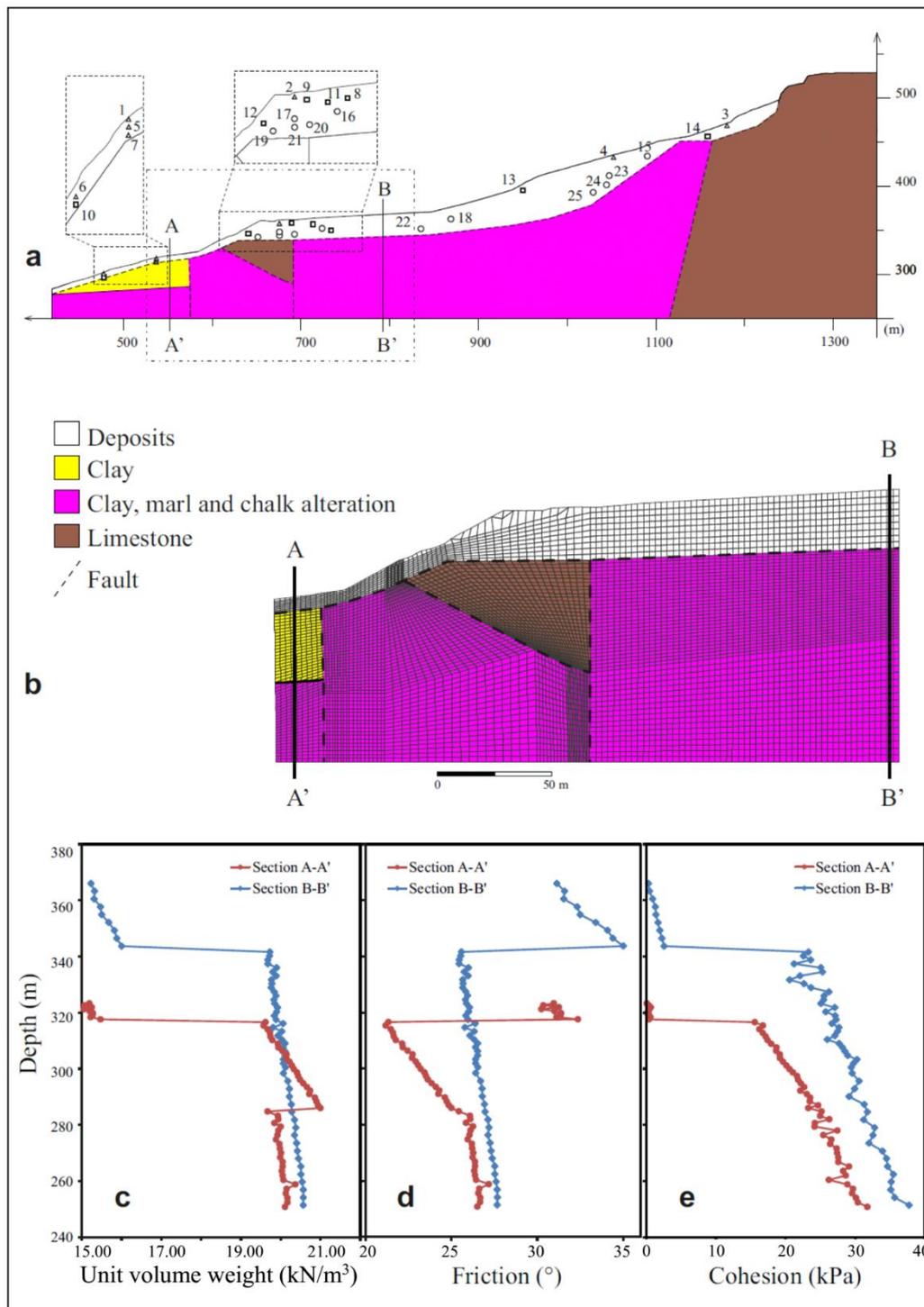
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349 **Fig. 6.** a) Geological section (section S-S', Fig. 5); b) grid used for modelling; c) density variation
 350 with depth; d) friction variation with the depth; e) bulk modulus and shear modulus variation with
 351 depth.



352
 353 A specific type of grid was selected (Fig. 6b), suitable for discretizing the soil with small mesh sizes
 354 and including existing discontinuities. The Mohr-Coulomb parameters for all materials are indicated
 355 in Table 1. Fixed mean values were used for limestone, while variable parameters were employed for
 356 other materials according to the above discussed equations, as we can see in Fig. 6c, d, e, where the

357 graphs show some parameter variations with depth along two verticals A-A' and B-B'(Fig. 6b). The
 358 Yielding Index F_{YI} was calculated in each grid cell belonging to deposits formation. In order to study
 359 the influence of the uncertain knowledge of the selected distribution of physical-mechanical
 360 parameters on the local material stability (and integrity) of slopes, the stress state distributions
 361 consequent to the 650 realizations were acquired.

362

363 **Table 1**

364 Mechanical parameters for each material.

365

	Unit volume weight (kN/m ³)	Friction (°)	Cohesion (kPa)	Elastic modulus (MPa)
Deposits	15.00 ÷ 17.00	30 ÷ 40	0 ÷ 5	3 ÷ 4
Clay	19.50 ÷ 21.00	21 ÷ 25	15 ÷ 25	10 ÷ 15
Clay, marl and chalk alteration	19.50 ÷ 21.50	25 ÷ 30	20 ÷ 50	25 ÷ 40
Limestone	24.00	40	50	100

366

367 At the end of each analysis, two particular FISH functions were employed, both related to each grid
 368 cell. The first, *ps3d.fis*, allowed calculating the principal stress and record them in an appropriate
 369 extra grid variables, the second, *mcfos.fis*, was used to calculate F_{YI} factor, based on the calculated
 370 values of the stress and on physical-mechanical parameters. However, in order to conveniently
 371 discuss the set of the results also from a statistical point of view, we grouped (by FISH and Fortran
 372 routines not reported) the whole ensemble of the F_{YI} – resulting from the 650 Monte Carlo multiple
 373 realizations and related to the 2943 meshes of the deposits formations – into different size unit
 374 intervals, called bins in the histogram rendering. In Fig. 7, the distribution of F_{YI} resulting from two
 375 different Monte Carlo realizations, with 0.25 size interval increments, are displayed. It is worth noting
 376 that the F_{YI} variability of certain areas is lower than the variability of other areas, as the more detailed
 377 discussion reported in the following section will confirm.

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386 **Fig. 7.** F_{YI} distribution in the deposits formations, 0.25 unit interval increments.



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4. Statistical analysis

390 The resulting F_{YI} numerical ensembles are analyzed by a statistical point of view. To this purpose,
391 some commonly utilized tools were selected and briefly described in Table A.1 (Appendix).

392
393

4.1. Global statistics insights

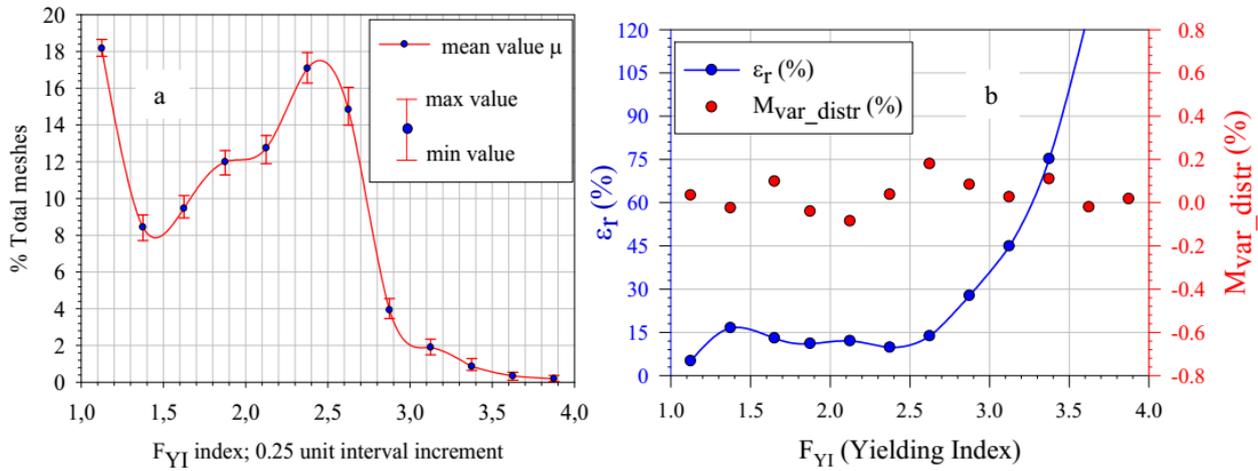
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395 For each of the 650 different realizations, the number of spatial meshes, each of which associated to
396 a small area, whose F_{YI} values were within a unit interval range around a selected mean, was counted.
397 Fig. 8a shows the resulting distributions and the variation of the percentage of the total meshes of the
398 deposits formations, grouped into 12 ensembles characterized by 0.25 unit interval increments
399 starting from $F_{YI}=1.125$. The inspection of Fig. 8a, suggested the introduction of two more
400 parameters. The first one was a Coefficient of Variability (COV, as named in statistics) and
401 accordingly defined as: $\epsilon_r = \Delta/\mu$, where Δ was the total variation of the percentages of the number
402 of meshes (the vertical bar length in Fig. 8a) counted for each of the 12 ensembles, while μ was the
403 mean value of the meshes number percentages counted over the 650 realizations. The other one was
404 the median difference defined as: $(M_{var_distr} - \mu)$, where $M_{var_distr} = (Max + Min)/2$, while
405 ‘Max’ and ‘Min’ were, respectively, the upper and the lower values of each bar in Fig. 8a.

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408

409 **Fig. 8.** a) Total distributions, 0.25 unit increment ($F_{YI} = 1.125$ as first value); b) Relative variability
 410 and median difference of meshes number within F_{YI} unit.

411



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414 Accordingly, the ε_r indicates how many areas, in percentage, changed their stress state when
 415 parameter values changed, in other words the sensitivity of areas stress state to the overall parameters
 416 variability, while M_{var_distr} is a heuristic indication of the skewness of the distribution of the number
 417 of meshes, characterized by a F_{YI} value. Thus, in Fig. 8b the relations between ε_r and M_{var_distr} vs
 418 F_{YI} are reported. It appears that the relative variability ε_r is correlated in some way to the F_{YI} index
 419 value, while the relation between the skewness and the index seems to be completely random.
 420 Observing the ε_r values, it appears that the relative variability of the meshes number, in some sense
 421 also connected to the resulting variability of areas characterized by a specific stress state, is much
 422 lower when the area stress state is close to critical conditions. Moreover, the trend of the ε_r parameter
 423 identifies three different F_{YI} fields: the first ranging from 1 to 1.45, the second from 1.45 to 2.5 and
 424 the last one from 2.5 to 4 (in the figure, F_{YI} values greater than 3.5 are out of the selected scale). In
 425 the first and last fields, where the stress state is respectively close to or far from critical conditions,
 426 the area extension affected by changing condition increases as the F_{YI} increases, while within the
 427 intermediate fields the relative extension variability is almost unaffected by the increase of Yielding
 428 Index. All this, at least, according to the selected mathematical models.

429

430 4.2. Local statistics insights

431

432 Previously, a global investigation was performed. An explorative analysis is performed about the
 433 issue of how the stress state related to some areas of the deposit layer was influenced by the uncertain
 434 knowledge of the numerical values of the most important mechanical parameters. To this end, 25

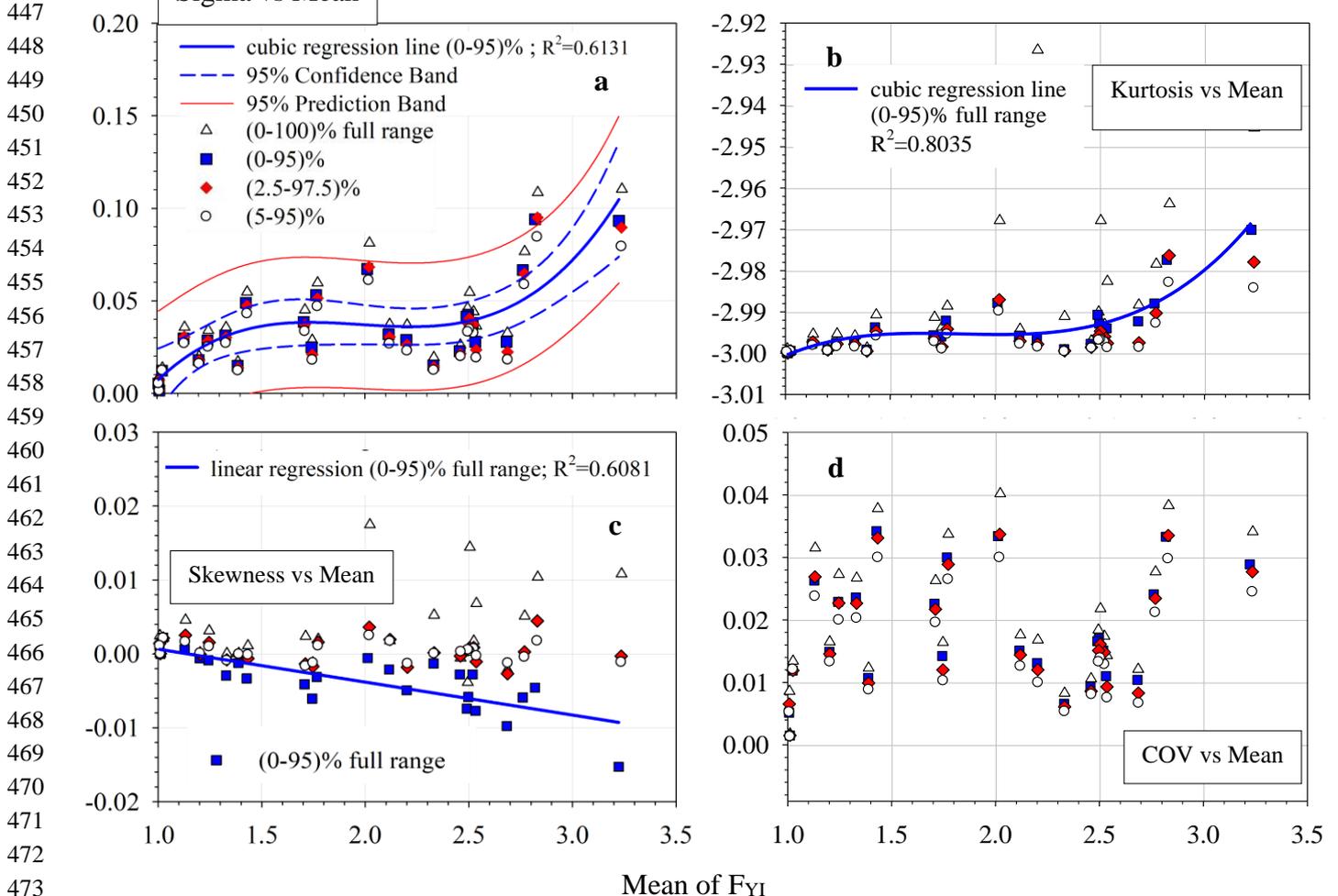
435 nodes, representative as much as possible of each critical region of the layer, labelled according to
 436 their progressive distance from the topographic line (depth) (Fig. 6a), were selected by inspecting
 437 Figs 6, 7. As far as possible, all meshes typologies and zones, through the superficial deposit, were
 438 identified and represented by one or more nodes.

439 An important explored parameter, required for comparing data, was the ‘Mean’ value of the statistical
 440 ensembles. In Fig. 9, the related comparisons are displayed. In this case, non-vanishing correlations
 441 seem to emerge between the Standard deviation (Fig. 9a, $R^2 = 0.6131$), the Kurtosis (Fig. 9b, Squared
 442 Pearson Correlation: $R^2 = 0.8035$), the Skewness (Fig. 9c, $R^2 = 0.6081$) and the Mean. On the other
 443 hand the COV parameter is clearly uncorrelated to the Mean (Fig. 9d).

444

445 **Fig. 9.** Sigma, Kurtosis, Skewness, COV vs mean.

446



474

475 The resulting fitting curves were selected according to the requirement of the best goodness of fit. In
 476 Fig. 9a, it is interesting to observe that the cubic regression line of the comparison of the Standard
 477 deviation $\sigma_{\Delta R2}$ (regarding only the F_{YI} values falling within the 0.% up to 95% percentile, Table A.1,

478 Appendix) to the mean $\mu_{\Delta R2}$ (values ranging within the interval 0-95%) is similar to the regression
 479 line of the parameter ε_r in Fig. 8b. Moreover, we grouped the 25 nodes, displayed in Fig. 6a, into
 480 three sets, characterized roughly by different depths and positions: superficial nodes numbered from
 481 1 to 7, depth ranging from 0.11 to 3.85 m; intermediate nodes from 8 to 14, depth from 4.16 to 6.81m;
 482 deep nodes from 15 to 25, depth from 11.90 to 35.53 m. In Table A3, the comparison between
 483 Standard deviations to Means, related to different percentiles range, in order to avoid outliers, are
 484 indicated. Also in this case, after many attempts aimed at identifying correlations trend, the 0%-95%
 485 percentiles range was selected to perform and discuss further analyses (in particular, $\mu_{\Delta R2}$ and $\sigma_{\Delta R2}$
 486 in bold characters in Table A3. Accordingly, hereafter only the 0%-95% percentiles range will be
 487 considered. The results are reported in Fig. 10.

488

489 **Fig. 10.** Sigma, Kurtosis, Skewness, COV vs Mean.

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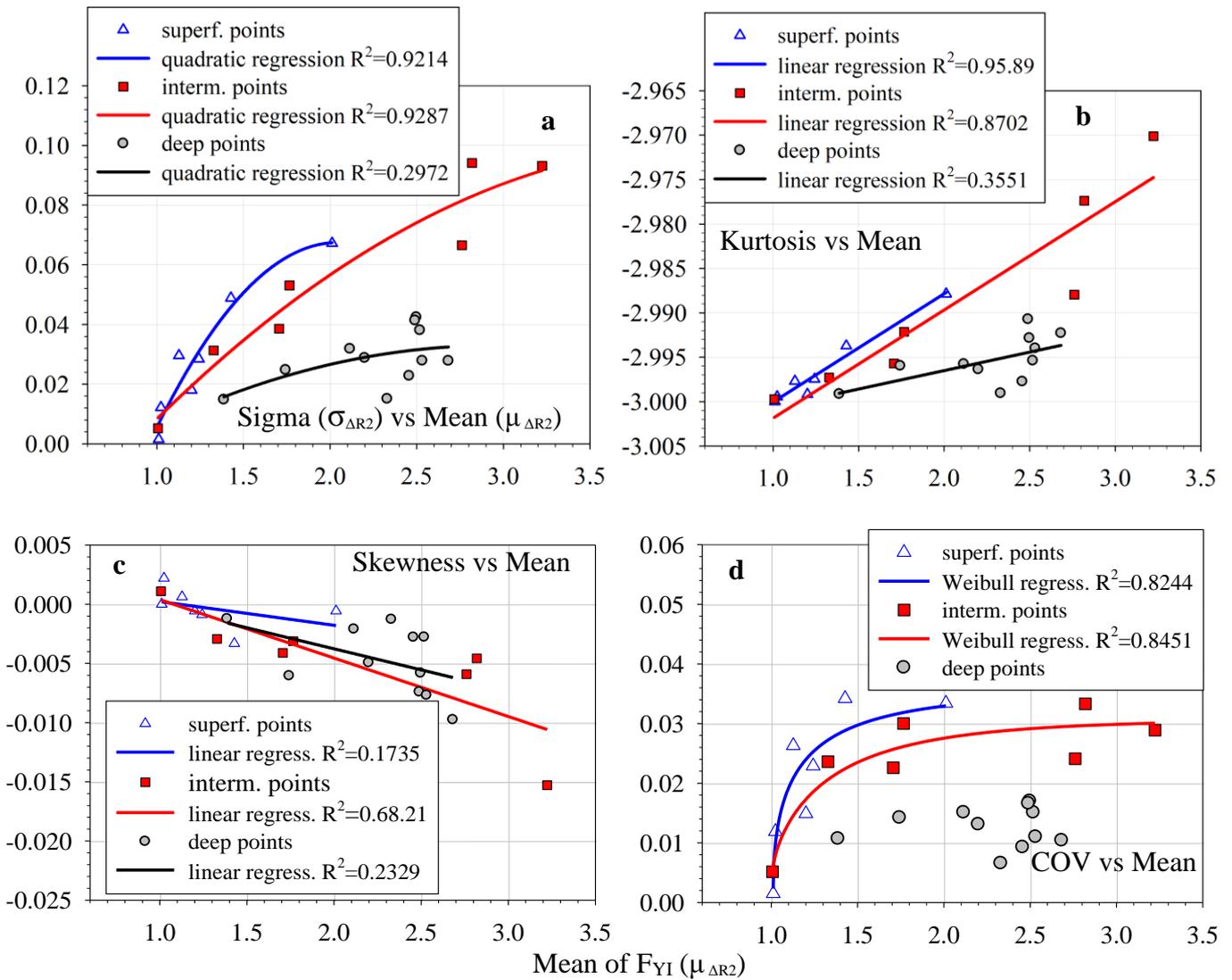
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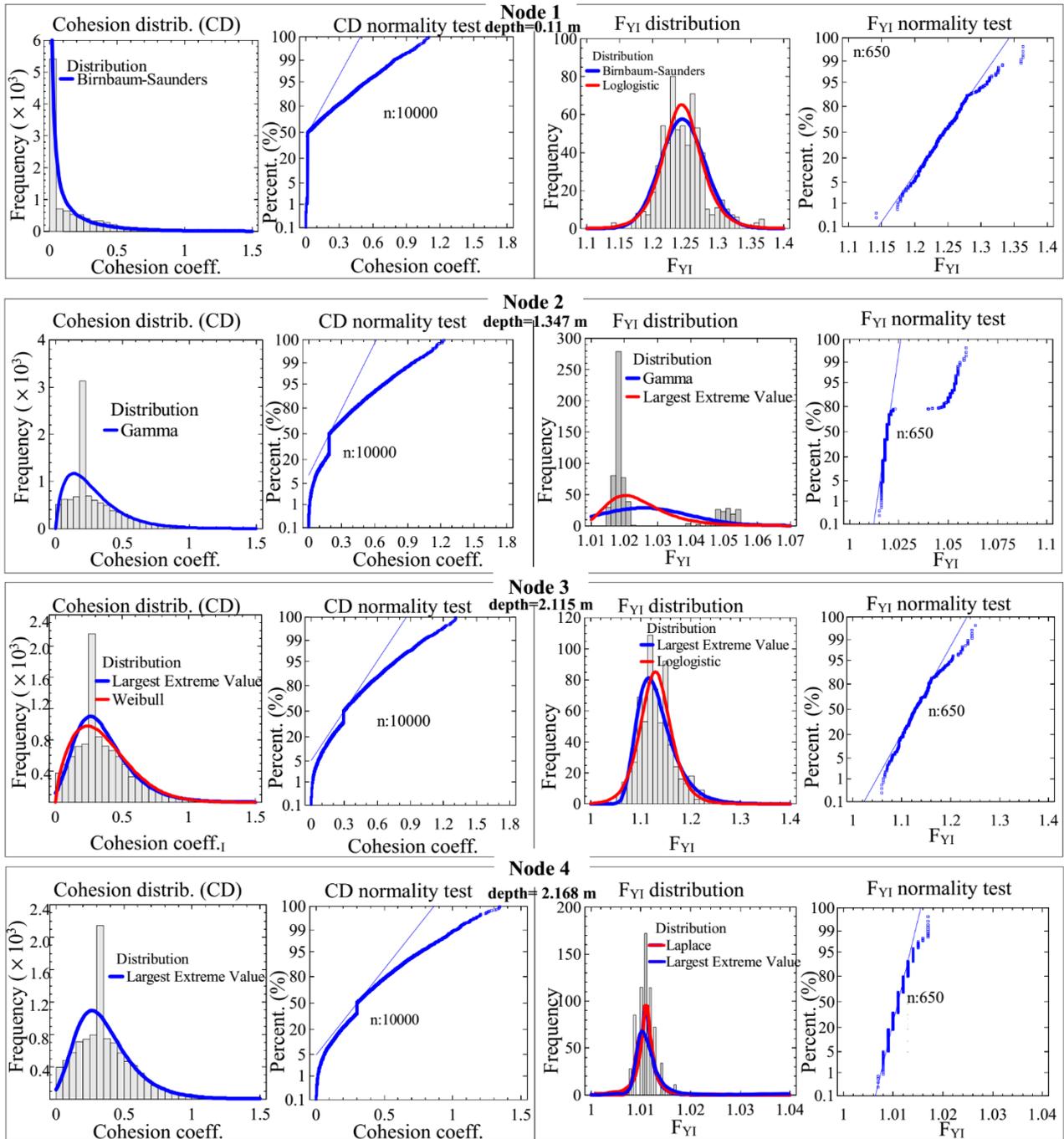
504 It appears clearly that for superficial and intermediate points, located within the first 7 m, strong
505 correlations between the Standard deviation (quadratic), Kurtosis (linear) and the Mean take place.
506 The related R^2 ranges between 0.8702 and 0.9287. The correlation between the COV and the Mean
507 is described by a Weibull fitting curve ($0.8244 < R^2 < 0.845$), indicating that just for low F_{YI} values
508 the two parameters are correlated. The Skewness is only weakly correlated to the Mean ($0.1735 < R^2$
509 < 0.6821). On the other hand, regarding nodes located in the third layer (depth higher than 7-12 m),
510 these kinds of correlations become low or are completely lost. A statistical computer code
511 (STATGRAPHICS Centurion XVI, 2010) was applied in order to analyse to what extent the initial
512 distributions of the numerical values, fitting the mechanical parameters input, are comparable with
513 the distributions of the numerical F_{YI} values resulting from the application of the mathematical model
514 discussed in the previous sections. The option called *Comparison of Alternative Distributions* on the
515 *Tables and Graphs* dialog box allows to select, among the 27 most common distributions, the best
516 curves fitting the data under study and to list them according to their goodness of fit, in accordance
517 with the ‘Maximum LogLikelihood Test’ (MLLT) (Cox and Hinkley, 1974, Casella and Berger,
518 2001). Tables A.4a and 4b report the results of the tests aimed at identifying the goodness of fit,
519 related to the statistical comparison of the cohesion coefficients distributions inputs and the resulting
520 F_{YI} distributions, for the selected nodes. The higher the value of the number along the columns, the
521 better the goodness of fit to data of the selected distributions.

522 For each node, bold blue and red colours indicate the numerical rank of the best fitting, respectively,
523 of the cohesion coefficients and the resulting F_{YI} selected distributions. In order to gain more insights
524 on the sensitivity of the resulting F_{YI} numerical values and of their distributions with respect to the
525 uncertainties associated with the cohesion coefficient (the selected parameter in order to study this
526 kind of issue), we detailed the performed analysis, reporting the related graphics of some selected
527 nodes, representative of particular areas of the slope under study (Figs. 11-15). In the figures, the blue
528 colour of the F_{YI} fitting curves indicates the same fitting distribution as that of the cohesion
529 coefficient, identified by MLLT, while the red colour indicates the best fitting curves of the F_{YI} values
530 ensemble, resulting from the 650 runs, identified by the same MLLT procedure. Since each
531 distribution was elaborated starting from the Standard Normal Distribution, a ‘Normality’ test was
532 also included to visualize how the geo-mechanical, the mathematical and the numerical (in particular
533 the discretization of the continuum through meshes) selected models could modify the final resulting
534 distributions. Moreover, it is worth noting that, according to the initial assumption that excludes (for
535 this step) any connectivity among the geo-mechanical parameters distributions of nodes, the
536 distributions of the cohesion coefficient of each node is completely independent from the others.

537

538

539 **Fig. 11.** Comparison between cohesion coefficients and F_{YI} fitting distributions related to superficial
540 nodes 1, 2, 3, 4.



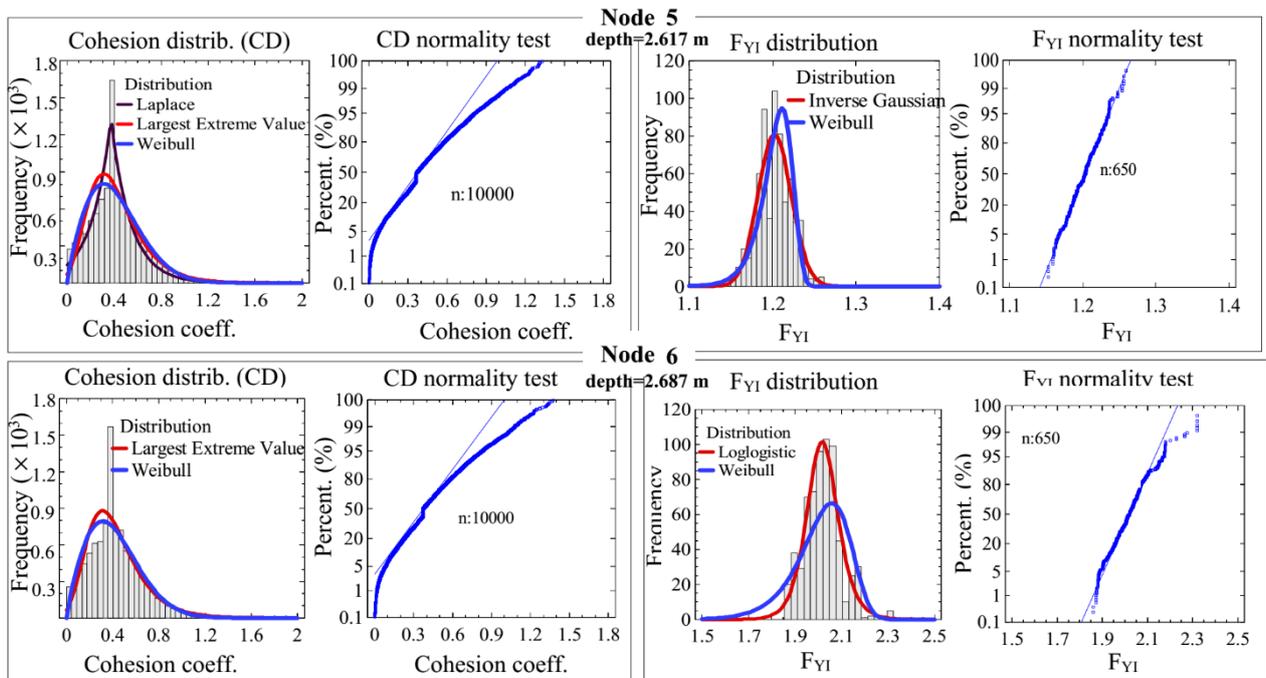
541

542 On the other hand, the distributions of the resulting F_{YI} coefficients may be affected not only by the
543 local features of the selected area, but also, and in a very complicate way, by the influence of the
544 entire system, at different scales. By consequence, the comparison of these two kinds of distributions
545 could be considered as an indication of how the uncertainties were transmitted (according to the
546 selected geological-mathematical and numerical models) through the system and which are the areas
547 more sensitive to the lack of knowledge of fundamental parameters, necessary to run the model. First

548 of all, we considered (Fig. 11) the superficial nodes 1, 2, 3 and 4, whose depths range from 0.11 m to
 549 2.168 m (Fig. 8a). The ‘Normality’ test failed for all the resulting distributions related to the
 550 implemented cohesion coefficient input distribution, indicating that – for these selected nodes – the
 551 application of the mathematical model strongly affected and perturbed the Standard Normal
 552 Distribution on which the Monte Carlo approach was based (Fig. 2b and c). In particular, for the
 553 Cohesion coefficient, the observed jump of the ‘Normality’ test curve and spike, if any, are a
 554 consequence of setting a random outlier value, resulting from Eq. (3) and (7), equal to the distribution
 555 mean (section 2). F_{YI} ‘Normality’ tests for nodes 1, 3 and 4 are verified for central data, but not for
 556 the tails. For node 2 the test failed completely. For these superficial nodes, the fitting F_{YI} distributions
 557 are completely different compared to the cohesion coefficient distribution. As regards these nodes, it
 558 could be inferred that the obtained F_{YI} distributions were strongly affected by non-local 'transmission'
 559 of the uncertainties that modified the local distribution of the assigned inputs parameters. In
 560 particular, node 1 is located at the toe of the slope (Fig. 8a), while node 2 (Fig. 8b) is located in an
 561 area that was discretized differently from the neighbouring regions. This occurrence could affect the
 562 outcomes of the models. Moreover, the resulting F_{YI} distribution of node 2, clearly indicates that the
 563 related area presents a criticality by a stability point of view.

564

565 **Fig. 12.** Comparison between cohesion coefficients and F_{YI} fitting distributions related to superficial
 566 nodes 5, 6.
 567



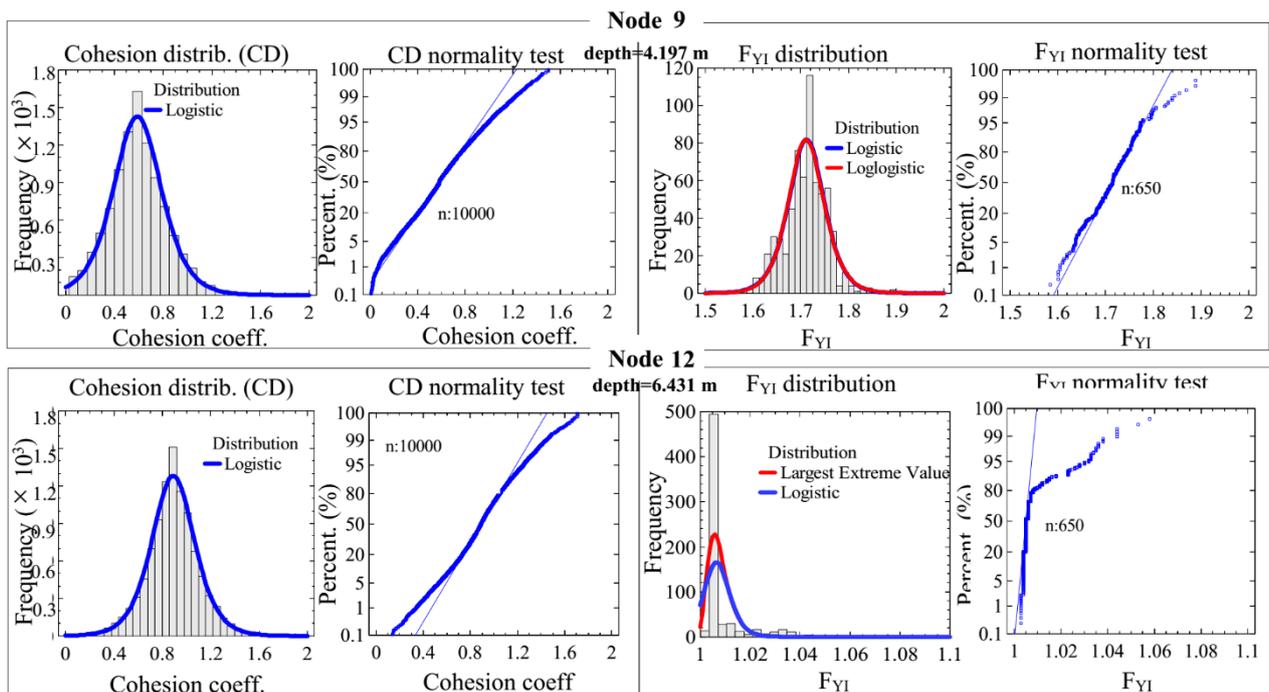
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570 Furthermore, the spreading of its distribution appears to be low. It is interesting that, despite the
 571 proximity of node 5 to node 1 (Fig. 6a) and the similarity of the related meshes (Fig. 6b), the two
 572 kinds of distributions and the resulting ‘Normality’ tests are different (Fig. 12). The only difference
 573 is the depth of their locations. Accordingly, it could be inferred that, at the toe of the slope, also a
 574 relative difference in the depth (about 2.5 m) could be significant. Node 6 is located at a different
 575 point of the toe, but at almost the same depth of the node 5 (2.687 m). For both nodes the cohesion
 576 coefficient distributions were the same (Weibull), but their F_{YI} distributions were different, indicating
 577 that there may be a different influence of the entire system on the two different meshes. Fig. 13 shows,
 578 in particular, the response of the intermediate node 12 that appears to be similar (except for the
 579 occurrence of a double mode distribution regarding the node 2) to both the distribution and the
 580 resulting ‘Normality’ test of node 2.

581

582 **Fig. 13.** Comparison between cohesion coefficients and F_{YI} fitting distributions related to
 583 intermediate nodes 9, 12.



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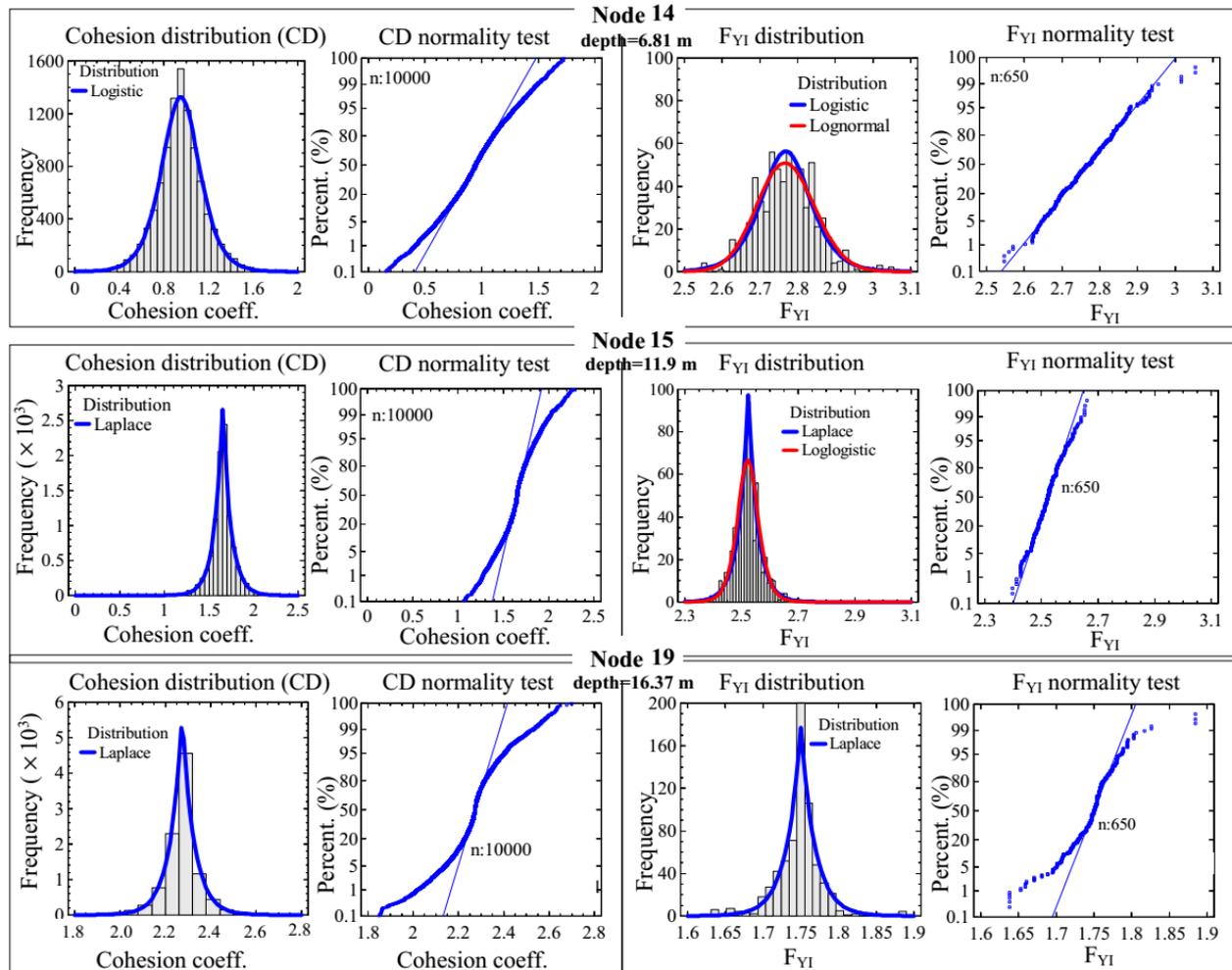
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586 Further, the node 12 experiences low values of both the mean and the spreading of its resulting F_{YI}
 587 distribution. It is interesting to note that both nodes 2 and 12 are located in an area that, reasonably,
 588 could be considered critical by a global stability of view. Nevertheless, also the mean and the
 589 spreading of node 4 are low. However, it does not seem to be located in a critical region of the slope
 590 (Fig. 6a). Accordingly, the particular response of the F_{YI} distribution could be ascribed to the
 591 occurrence of just local instability, not necessarily due to an actual phenomena, but also to the applied

592 model. Moreover, node 9, although located in an area close to node 2, thus with the same type of
 593 discretization, did not experience the same F_{YI} distribution characteristics. We can deduce that
 594 stability conditions, in this case, are more influential than the type of the selected mesh. Finally, in
 595 Figs 14 and 15, the responses of an intermediate node (14) and deep nodes (15, 19, 23, 24, 25).

596

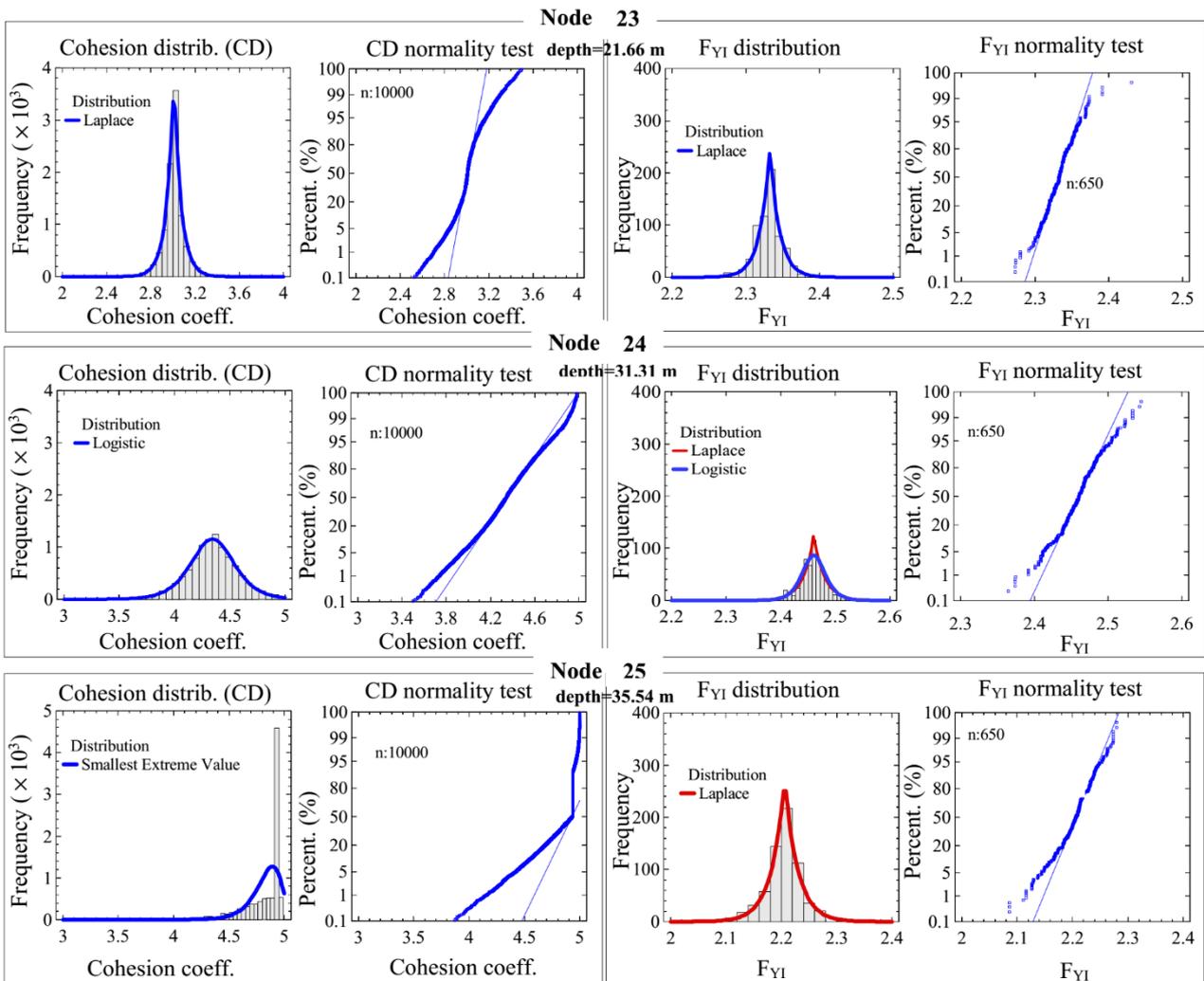
597 **Fig. 14.** Comparison between cohesion coefficients and F_{YI} fitting distributions related to
 598 intermediate node 14 and deep nodes 15, 19.



599

600 Based on the previous discussion, since the F_{YI} distribution is very similar to the correspondent
 601 cohesion coefficient, all these nodes, except for the last one, node 25, appear to be weakly influenced
 602 by the global response of the system. Node 25 was placed at the point where the distance from the
 603 corresponding topographic is maximum for the selected region of the slope. Therefore, similarly to
 604 what happens for node 1, the particular selected mathematical model (based on a linear increase of
 605 the mechanical coefficients as the depth increases, eq. 6) strongly affected the resulting cohesion
 606 coefficient. However, it is interesting to note that the resulting F_{YI} distribution (also for node 1, but
 607 in particular for node 25) is completely different. In addition, by observing the answer of other similar
 608 points, 23, 24, 25, this occurrence suggests that, in this case, the spatial scale of the 'influence area'

609 (in which the mechanical response of each node is related to that of the other nodes) is medium, that
 610 is, non-local (single mesh), non-global (the whole system) but includes a limited number of meshes.
 611
 612 **Fig. 15.** Comparison between cohesion coefficients and F_{YI} fitting distributions related to deep nodes
 613 23, 24, 25.



614
 615 In summary, observing the previous figures (Figs. 11-15) and in particular Table A.4a and b fitting
 616 ranks, we can infer that as regards superficial nodes (0.11m - 3.85m; 1-7 nodes) the resulting F_{YI}
 617 distributions is different to the cohesion coefficient distributions (the blue and red fitting ranks are
 618 not within the same columns under the same fitting distribution), except for node 3, evidencing that
 619 the mechanical local response (single mesh) is strongly influenced by different spatial scales (medium
 620 up to large, thus more than few meshes). The same applies to intermediate nodes (4.16m - 6.81m; 8
 621 - 14 nodes), except node 9. On the other hand, with respect to deep nodes (11.90m - 35.54m; 15-25
 622 nodes), similar distributions, Laplace fitting curve, are observed for seven nodes with the exception
 623 of 4 nodes (15, 20, 24, 25).

624

625 **5. Conclusions**

626 In this paper, the statistics of the numerical stress states resulting from different realizations of as
627 realistic as possible mathematical model of an actual slope, characterized by heterogeneity at multiple
628 scales, is presented and discussed, in order to study uncertainties propagation within the selected
629 system. To make the analyses as simple as possible, we considered the selected test case in a dry
630 condition, avoiding considerations about the saturation state of the system. The spatial distribution of
631 physical-mechanical parameters was realized by random Monte Carlo assignment to each point of a
632 numerical value belonging to a probabilistic distribution ensemble, assumed valid for the system
633 under study, added to a trend along the depth, identified by mechanical, physical and geotechnical
634 considerations.

635 Compared to previous approaches adopted also by us, instead of identifying a global and unique
636 safety factor through degradation of mechanical parameters until rupture, in this paper we focused on
637 the Yielding Index F_{YI} , which can serve as a proxy for identifying a local stress state, approaching
638 unitary value when close to local rupture conditions. This type of index can allow the identification
639 of areas particularly susceptible to becoming locally unstable and which may also induce global
640 instabilities. The numerical evaluation of this parameter depends on how the numerical model is
641 performed, including also the treatment of the intrinsic uncertainties. However, the initial local
642 assignment of a selected uncertainty could be heavily affected by the mechanical influence of the
643 entire system. Accordingly, in this paper we proposed an approach to analyze how the variability and
644 uncertainties propagate within the system under study, in order to identify zones where the
645 uncertainties related to instability parameters may grow beyond the initial assignment.

646 The proposed algorithms were successfully implemented in a commercial code (FLAC2D) through
647 an interactive tool (FISH) and then applied to an actual site as test case, with particularly complex
648 lithostratigraphic characteristics. The discussed procedure and results highlighted the effectiveness
649 and flexibility of the implementation. The analysis of the resulting statistical ensemble of the F_{YI}
650 outcomes, related to twenty five nodes located in different places of the test slope and grouped into
651 three levels, selected on the basis of their distances from the vertical point located on the topographic,
652 seems to reveal, for superficial and intermediate levels (0.11m - 7m), an interesting correlation
653 between the spreading of such parameters and their mean values (Fig. 10a). In particular, it appears
654 that the variability, indicated in Fig. 10a as Sigma $\sigma_{\Delta R2}$ of the calculated F_{YI} , decreases as the value
655 of the F_{YI} mean, obtained for each mesh by the calculations discussed above (Mean $\mu_{\Delta R2}$ in Fig. 10a),
656 decreases. Thus, it appears that this kind of analysis indicates lack sensitivity to overall variability of
657 the parameters affecting the mechanical response of superficial and intermediate layers zones (Fig.
658 10a) whose stress state is close to local stability criticality, at least from the particular typology of the

659 test case and the selected models. Namely, for the selected mathematical model and test case, when
660 the local area under study is close to critical conditions, thus with low F_{YI} values, the probable
661 statistical stress state distributions, resulting from multiple realizations of realistic set of numerical
662 values of parameters, is less sensitive to the randomness of the available information; whereas, when
663 the area is far from criticality, thus more stable, its stress state is more dependent on randomness. In
664 a thermodynamics and entropic sense, it may be argued that since close to criticality just few states
665 are accessible (low variability of the calculated F_{YI} coefficients), accordingly, the system is more
666 'ordinate', by a statistical point of view, than in stable conditions in which the number of ways to be
667 stable is higher (high F_{YI} variability). Moreover, the comparison between the resulting F_{YI} distribution
668 and the cohesion coefficient input, selected as an important mechanical parameter to which
669 performing the comparison, suggests that some areas of the slope may be more influenced than others
670 by what happens in more or less far areas, identifying a spatial connectivity and uncertainties
671 transmission through local, medium and global scale of influence (Figs. 11-15).
672 The test slope system can be considered as common in engineering geology practice and, therefore,
673 the proposed methodology can provide useful indications for similar or also different situations, in
674 identifying areas of actual systems where any additional experimental surveys and more accurate
675 numerical modelling could be focused, excluding less important zones by the above discussed point
676 of view, saving time and economical efforts, important issues even for practitioners.

677

678 **Acknowledgment**

679 The authors would like to thank three anonymous referees and the Editor for the suggestions that
680 helped to improve the clarity and the quality of the paper.

681

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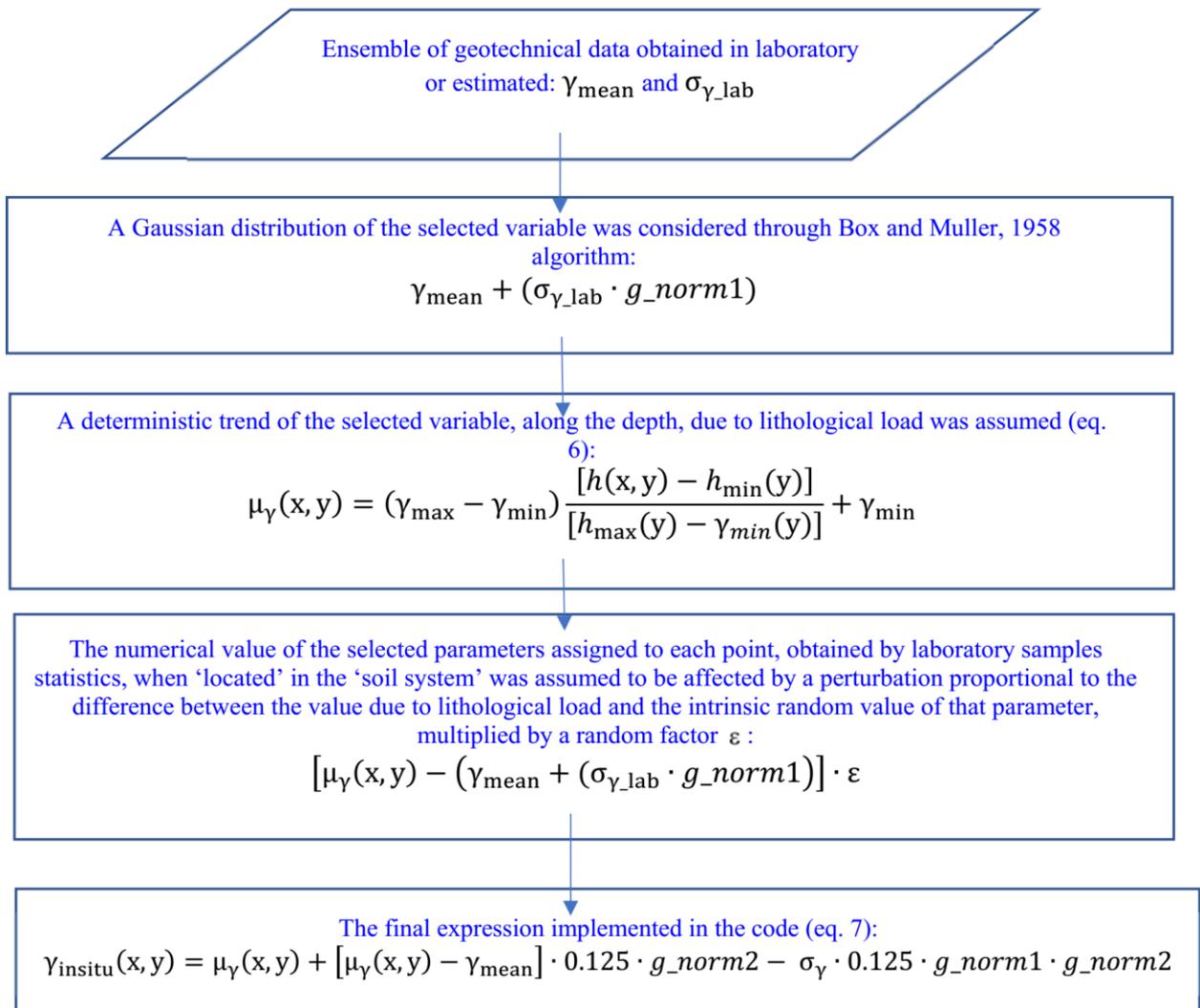
828 **Appendix**

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830 **Fig. A.1.** Flow Chart of the Numerical modelling

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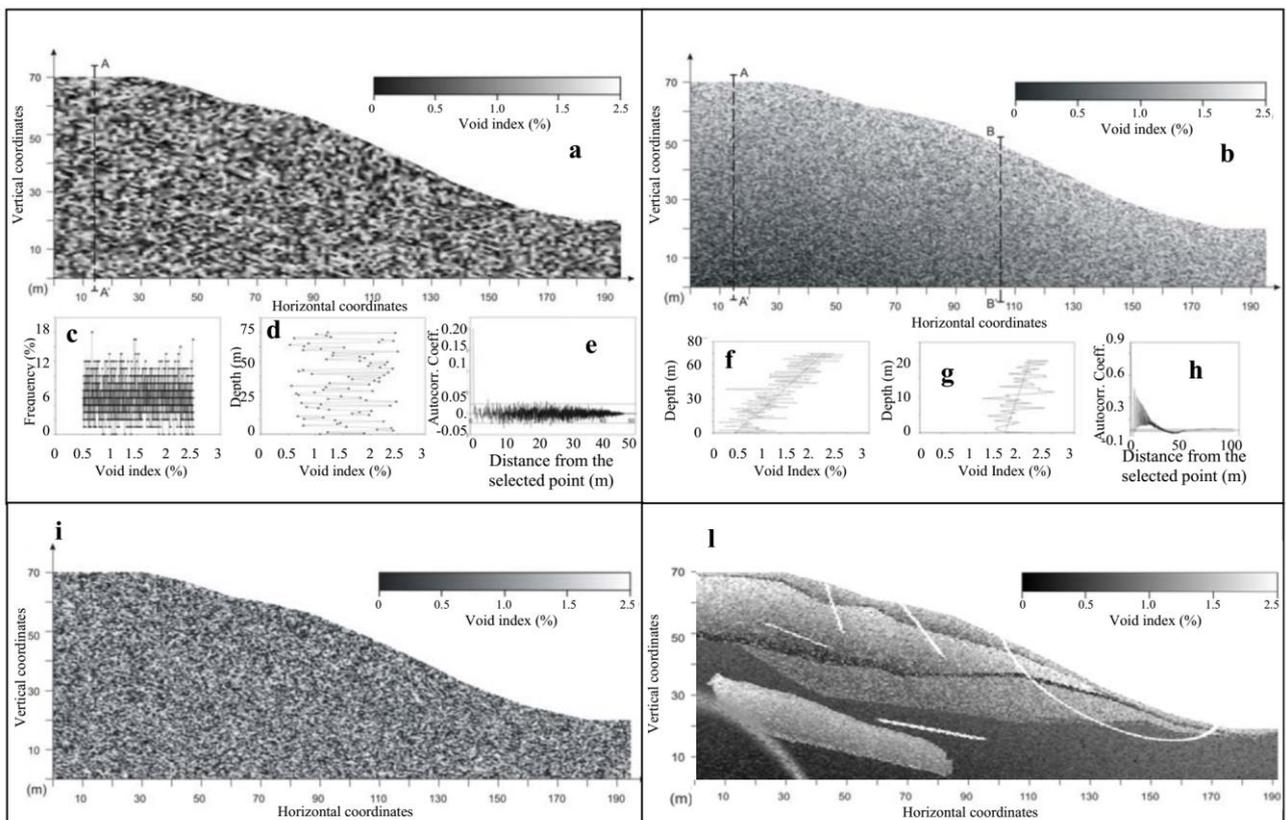
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843 **Fig. A.2.** Explanatory figures of the adopted Monte Carlo model (eq.7): **a)** Numerical realization,
 844 through 7437 nodes, without a deterministic trend of the void distribution, regarding a synthetic2D
 845 slope section; **b)** Numerical realization, through 31408 nodes, with the lithological load of the void
 846 distribution; **c)** Uniform pseudo-random realization used for the Monte Carlo application; **d)** Void
 847 distribution along the section A-A'; **e)** Autocorrelation between a selected node and all the other
 848 nodes, showing no correlation (inside Bartlett's limit, Snedecor and Cochran, 1989
 849 $|Autocorrelation| \leq \pm 1.96\sqrt{N}$ where N is the number of the considered points); **f)** Void distribution
 850 along the section A-A' related to the Fig. A.2.b; **g)** Void distribution along the section B-B' related
 851 to the Fig. A.2.b; **e)** Autocorrelation between a selected node and all the other nodes, showing
 852 correlation (outside Bartlett's limit) due to lithological applied trend; **i)** Numerical realization,
 853 through 31408 nodes, without a deterministic trend of the void distribution; **l)** Realization of a
 854 synthetic multilayer 2D section, characterized by faults and weakness lines.

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Fig. A.3a Distance from topographic FISH routine.

```
def distance
  var1=0.0
  var2=10000
  ;-----
  loop i (1, izones)
  ;-----
    loop j (1, jzones)
      yc = (y(i,j)+(y(i+1,j)+ y(i,j+1)+(y(i+1,j+1)))/4.0
      ;yc = approximate centroid of the selected zone
      xx = x(i,j)
      yy = table(1,xx)
      zc =yy - yc
      ;zc = h(x,y) in Fig. 1
      hmax = max(zc, var1)
      ;hmax = hmax(x,y) in Fig. 1
      hmin = min(zc, var2)
      ;hmin = hmin(x,y) in Fig. 1
      var1=hmax
      var2=hmin
    end loop
  ;-----
end loop
;-----
end
distance
```

Fig. A.3b Random density value realization FISH routine.

```
def property
;-----
loop i (1, izones)
;-----
  loop j (1, jzones)
    rand1 = urand
    rand2 = urand
    rand3 = urand
    rand4 = urand
    dmax = 1700
    dmin = 1500
    dmed = (dmax+ dmin)/2
    diff = dmax - dmin
    yc = (y(i,j)+(y(i+1,j)+ y(i,j+1)+(y(i+1,j+1)))/4.0
    xx = x(i,j)
    yy = table(1,xx)
    zc =yy-yc
    g_norm1=(-2.*ln(rand1))^0.5*cos(2.*pi*rand2)
    g_norm2=(-2.*ln(rand3))^0.5*cos(2.*pi*rand4)
    ;pi = p greek
    ;g_norm1, g_norm2: Box, Muller; Eq. (3)
    mu=((dmax-dmin)/(hmax-hmin))*(zc-hmin)+dmin
    ;mu = Eq.(6)
    phys=((dmin+dmax)/2+((dmax-dmin)/8.)*g_norm1
    ; phys = Eq.(1) and Eq.(2)
    density(i,j)=phys+[mu-phys]*[1+0.125*g_norm2]
    ; density(i,j)=Eq.(5), Eq.(6) and Eq.(7)
  ;-----
  if density(i,j)<dmin then
    density(i,j)=mu
  end if
  ;-----
  if density(i,j)>dmax then
    density(i,j)=mu
  end if
  ;-----
  end loop
;-----
end loop
;-----
end
property
```

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874 **Table A.1** Exploited statistical tools

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Name	Mathematical expressions	Notes
n		Number of realizations (650)
median		F _{YI} value corresponding to 50% percentage of the cumulative distribution function
Cumul(·)		F _{YI} value corresponding to the percentage equal to the value included in the brackets (·) of the cumulative distribution function
Ensembles embedding only the F _{YI} values falling within the following ranges: $\Delta R1=$ $\Delta R2=$ $\Delta R3=$ $\Delta R4=$	$[\text{Cumul}(100\%), \text{Cumul}(0.0\%)]$ $[\text{Cumul}(95.0\%), \text{Cumul}(0.0\%)]$ $[\text{Cumul}(97.5\%), \text{Cumul}(2.5\%)]$ $[\text{Cumul}(95.0\%), \text{Cumul}(2.5\%)]$	Suffixes used in Table A.3
Sample moment of <i>r</i> -th order	$m_r = \frac{1}{n} \sum_{i=0}^n (x_i - \hat{\mu})^r$	Statistical central moment about the mean: $\mu = \frac{1}{650} \sum_{i=1}^{650} (F_{YI})_i$
Standard Deviation	$\hat{\sigma} \equiv m_2^{1/2} = \sqrt{\frac{1}{n} \sum_{i=1}^n (x_i - \hat{\mu})^2} \equiv \sqrt{\frac{1}{650} \sum_{i=1}^{650} [(F_{YI})_i - \hat{\mu}]^2}$	Provides a measure of the data spread around the sample mean
Skewness	$\hat{\gamma}_1 \equiv \frac{m_3}{\hat{\sigma}^3} = \frac{\left[\frac{1}{n} \sum_{i=1}^n (x_i - \hat{\mu})^3 \right]}{\sqrt{\left[\frac{1}{n} \sum_{i=1}^n (x_i - \hat{\mu})^2 \right]^3}}$	Asymmetry of the data around the median.; a negative skew (left) has fewer large values and a positive skew (right) has fewer low values.
Kurtosis	$\hat{k} \equiv \frac{m_4}{\hat{\sigma}^4} = \frac{\left[\frac{1}{n} \sum_{i=1}^n (x_i - \hat{\mu})^4 \right]}{\sqrt{\left[\frac{1}{n} \sum_{i=1}^n (x_i - \hat{\mu})^2 \right]^4}} - 3$	A high kurtosis value indicates that the distribution has a sharp peak with long and fat tails.
COV	$= \hat{\sigma} / \mu$	Coefficient of variation, indicating the dispersion of the distribution, normalized to the mean

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890 **Table A.2** Fitting curves.

Distribution Name	Distribution fitting function		Notes
Birnbaum-Saunders	$a = \sqrt{(x_i - \alpha)/\beta}$ $z = (a^2 - 1)/a\gamma$ $\varphi = \text{Nor. Stand. Distr.}$	$[(a^2 + 1)/2a\gamma(x_i - \alpha)] \cdot \varphi(z)$	Fitting parameters: α = location β = scale γ = shape
Gamma	$\Gamma(c) = \int_0^\infty t^{c-1} e^{-t} dt$ Gamma function (Integral form)	$\{1/[b \cdot \Gamma(c)]\}(x_i/b)^{(c-1)} \exp(-x_i/b)$	Fitting parameters: c = shape b = scale
Inverse gaussian	$z_i = \ln(x_i/\theta)$ $\varphi = \text{Nor. Stand. Distr.}$	$\{\sqrt{\beta}/[x_i \cdot \exp(z_i/2)]\} \cdot \varphi(z_i)$	Fitting parameters: θ = mean β = scale
Laplace		$(\lambda/2) \cdot \exp[-\lambda x - \mu]$	Fitting parameters: μ = mean λ = scale
Largest and Smallest Extreme Value	$z_i = (x_i - \alpha)/\beta$	$\{\exp[-z_i - \exp(z_i)]\}/\beta$	Fitting parameters: α = location β = scale
Logistic	$z_i = (x_i - \mu)/\sigma$	$\exp(z_i) / \{(1/\sigma) \cdot [1 + \exp(z_i)]^2\}$	Fitting parameters: μ = mean σ = standard deviation
Loglogistic	$z_i = (\ln x_i - \mu)/\sigma$	$\exp(z_i) / \{[1/(\sigma \cdot x_i)] \cdot [1 + \exp(z_i)]^2\}$	Fitting parameters: $\exp(\mu)$ = median σ = scale
Lognormal		$[1/(\sqrt{2\pi} \cdot \sigma \cdot x_i)] \exp[-(\ln x_i - \mu)^2/(2\sigma^2)]$	Fitting parameters: μ = mean σ^2 = variance
Normal		$[1/(\sqrt{2\pi} \cdot \sigma)] \exp[-(x_i - \mu)^2/(2\sigma^2)]$	Fitting parameters: μ = mean σ^2 = variance
Weibull		$(\alpha/\beta^\alpha) \cdot x_i^{(\alpha-1)} \cdot \exp[-(x_i/\beta)^\alpha]$	Fitting parameters: α = shape β = scale

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903 **Table A.3** Results of the numerical experiments.

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mesh numb	depth (m)	level	$\mu_{\Delta R1}$	$\mu_{\Delta R2}$	$\mu_{\Delta R3}$	$\mu_{\Delta R4}$	$\sigma_{\Delta R1}$	$\sigma_{\Delta R2}$	$\sigma_{\Delta R3}$	$\sigma_{\Delta R4}$
1	0.11	superf	1.2468	1.2424	1.2461	1.2457	0.03398	0.02846	0.02827	0.02508
2	1.35	superf	1.0258	1.0243	1.0253	1.0247	0.0137	0.0122	0.0123	0.0124
3	2.12	superf	1.1324	1.1276	1.1313	1.1307	0.0357	0.0296	0.0305	0.0270
4	2.17	superf	1.0111	1.0109	1.0111	1.0111	0.0018	0.00152	0.00153	0.00137
5	2.62	superf	1.2027	1.2006	1.2026	1.2027	0.0199	0.0180	0.01754	0.0160
6	2.69	superf	2.0224	2.0118	2.0197	2.0191	0.0812	0.0672	0.0682	0.0611
7	3.85	superf	1.4337	1.4276	1.4332	1.4334	0.0548	0.0489	0.0475	0.0431
8	4.16	interm	2.8327	2.8193	2.8308	2.8299	0.1086	0.0941	0.0949	0.08455
9	4.20	interm	1.7118	1.7063	1.7110	1.7112	0.04512	0.0386	0.03722	0.03356
10	4.69	interm	1.7729	1.7661	1.7726	1.7722	0.0597	0.0531	0.0514	0.0470
11	6.33	interm	1.3319	1.3278	1.3316	1.3317	0.0356	0.0313	0.0302	0.0271
12	6.43	interm.	1.0087	1.0071	1.0079	1.0073	0.0087	0.0052	0.0066	0.0053
13	6.66	interm	3.2376	3.2236	3.2357	3.2355	0.1104	0.0932	0.0896	0.0794
14	6.81	interm	2.7697	2.7606	2.7686	2.7684	0.07665	0.06657	0.06485	0.05878
15	11.90	deep	2.5239	2.5184	2.5234	2.5233	0.04412	0.03810	0.03727	0.03262
16	12.13	deep	2.5051	2.4976	2.5037	2.5035	0.0546	0.04247	0.04049	0.03467
17	12.98	deep	2.1194	2.1147	2.1192	2.1189	0.03726	0.03185	0.03053	0.0268
18	13.07	deep	2.4951	2.4900	2.4958	2.4957	0.04604	0.04137	0.03781	0.03324
19	16.37	deep	1.7468	1.7436	1.7472	1.7475	0.02871	0.02479	0.02099	0.01806
20	16.82	deep	1.3894	1.3873	1.3893	1.3893	0.0171	0.01483	0.01375	0.01229
21	17.38	deep	2.5366	2.5318	2.5362	2.5364	0.03634	0.02785	0.02353	0.0192
22	18.63	deep	2.6866	2.6829	2.6871	2.6877	0.0326	0.02784	0.02232	0.01813
23	21.66	deep	2.3324	2.3300	2.3321	2.3320	0.01932	0.0151	0.0142	0.01249
24	31.31	deep	2.4598	2.4567	2.4598	2.4612	0.02612	0.0228	0.02131	0.01994
25	35.54	deep	2.2037	2.1998	2.2035	2.2039	0.03696	0.0288	0.02635	0.02296

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Table A.4a Goodness of fit by the MLLT.

Node numb	Depth (m)	Birnbaum-Saunders		Gamma		Inverse Gaussian		Laplace		Larg. Extr. (L)	Smal. Val. (S)
		Cohes.	F _{YI}	Cohes.	F _{YI}	Cohes.	F _{YI}	Cohes.	F _{YI}		
1	0.11	12115	1280	10372	1279	11948	1280	3722	1267	5196 L	1254 L
2	1.35	no fit	1872	3729	1870	no fit	1872	2787	2032	3657 L	2050 L
3	2.12	no fit	1251	1891	1249	no fit	1251	1643	1232	2136 L	1262 L
4	2.17	-1548	3181	1847	3181	-2440	3181	1649	3196	2085 L	3184 L
5	2.62	-2940	1623	833	1623	no fit	1623	934	1592	1189 L	1582 L
6	2.69	no fit	718	796	716	no fit	718	863	715	1145 L	720 L
7	3.85	-4062	965	-358	965	no fit	965	172	940	-3 L	923 L
8	4.16	no fit	523	-257	522	no fit	523	356	511	10 L	501 L
9	4.20	-3827	1093	-409	1093	no fit	1093	189	1099	-143 L	1048 L
10	4.69	-2960	910	-313	910	no fit	910	223	877	-241 L	865 L
11	6.33	-1456	1245	370	1245	no fit	1245	1261	1219	-159 L	1189 L
12	6.43	-391	2168	720	2164	no fit	2167	1364	2472	117 L	2489 L
13	6.66	-9971	512	781	511	no fit	512	1588	526	153 L	459 L
14	6.81	-486	749	1022	748	no fit	749	1701	721	303 L	709 L
15	11.90	5807	1107	5897	1107	5807	1107	7063	1108	no fit	1061 L
16	12.13	6362	975	6448	973	6362	975	7656	1039	no fit	944 L
17	12.98	7119	1217	7180	1217	7119	1217	8561	1242	no fit	1154 L
18	13.07	7273	1075	7322	1077	7273	1075	8673	1111	no fit	1042 S
19	16.37	10455	1384	10516	1384	10455	1384	12397	1476	no fit	1249 L
20	16.82	10969	1723	10998	1723	10969	1723	12816	1729	8717 S	1644 L
21	17.38	11110	1236	11135	1235	11110	1236	12965	1384	no fit	1142 L
22	18.63	11322	1302	11340	1303	11322	1302	13089	1449	8777 S	1120 S
23	21.66	8683	1648	8704	1646	8683	1648	10377	1712	no fit	1617 L
24	31.31	280	1446	314	1446	280	1446	285	1467	320 S	1380 S
25	35.54	2091	1236	2226	1231	2091	1236	3803	1327	5250 S	1193 L

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Table A.4b Goodness of fit by the MLLT.

Node numb	Depth (m)	Logistic		Loglogistic		Lognormal		Normal		Weibull	
		Cohes.	F _{YI}	Cohes.	F _{YI}	Cohes.	F _{YI}	Cohes.	F _{YI}	Cohes.	F _{YI}
1	0.11	3056	1286	10885	1288	11400	1280	2347	1276	10613	1176
2	1.35	2408	1882	3493	1888	2618	1872	1912	1866	3716	1528
3	2.12	1525	1250	1452	1253	438	1251	1270	1244	2091	1142
4	2.17	1519	3187	1381	3187	467	3181	1265	3180	2060	2815
5	2.62	968	1614	322	1614	-722	1623	856	1622	1217	1514
6	2.69	886	724	321	728	-773	718	751	710	1150	585
7	3.85	303	960	-585	960	-1844	965	230	965	194	913
8	4.16	459	518	-377	520	-1712	523	376	521	293	455
9	4.20	323	1102	-543	1102	-1800	1093	243	1092	153	997
10	4.69	367	904	-329	905	-1439	910	272	909	164	851
11	6.33	1288	1243	686	1243	-511	1245	1088	1245	884	1192
12	6.43	1403	2308	955	2314	93	2167	1214	2159	1004	no fit
13	6.66	1590	526	1135	527	no fit	512	1334	510	1038	417
14	6.81	1727	748	1343	749	372	749	1514	747	1242	666
15	11.90	6589	1113	6523	1113	5820	1107	5975	1106	4845	921
16	12.13	7184	1018	7109	1020	6374	975	6530	968	no fit	no fit
17	12.98	7954	1233	7905	1234	7127	1217	7235	1216	5900	993
18	13.07	8109	1097	8073	1097	7282	1076	7351	1079	5594	1006
19	16.37	11638	1441	11609	1439	10464	1384	10584	1386	no fit	662
20	16.82	12025	1737	12006	1737	10972	1723	11029	1723	9196	1497
21	17.38	12177	1331	12164	1331	11114	1236	11159	1232	no fit	no fit
22	18.63	12337	1392	12325	1390	11325	1302	11356	1304	9278	no fit
23	21.66	9643	1707	9628	1708	8686	1648	8723	1643	no fit	no fit
24	31.31	433	1465	398	1465	282	1446	360	1447	-96	1284
25	35.54	3188	1316	2937	1316	2937	1236	2480	1221	5040	no fit

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